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James,

George Oscar

1899

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On the Differential Equations  
Connected with Hypersurfaces

Dissertation

Presented to the Board of University  
Studies of the Johns Hopkins  
University for the Degree of  
Doctor of Philosophy

by  
George Oscar James

Baltimore

1899







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To Professor Ames I am grateful for the kind way in which he has helped me in my work in physics. Here his advice has been invaluable both during my undergraduate and graduate courses.



In the course of this work certain partial differential equations of special forms have arisen and it is these which underlie first the study and development of the theory of hypersurfaces.

The generalized Poisson equations and the equations of the type of Laplace in three independent variables are intimately connected with the existence of a hypersurface corresponding to two given fundamental forms, and the generalization of the notion of conjugate lines and lines of curvature leads to important information concerning these latter.

Later I hope I again take up the study of these equations.





# I

If we define a curved space of  $n-1$  dimensions as a manifoldness of  $n-1$  dimensions, which is contained in homaloidal (Euclidian) space of  $n$  dimensions and which can be moved about in this  $n$ -dimensional space without alteration of its linear element, then the coordinates of a point on the curved space can be represented by  $n$  variables and a single equation connecting them will serve to determine the curved space. Again,  $n$  such equations will determine a curved space of  $n-m$  dimensions  $n-2, 2$  surface and



is a twisted curve.

The theory of curves of multiple curvatures in  $n$ -dimensional space has been developed to some extent and very interesting results appear in two papers by Brunel\* and Poincaré\*\*. In his thesis presented to this University in 1898 M. R. 40-8936 found the Serret-Frenet formulae for curves of triple curvature by the binormal method of Darboux, and in a paper in the American Journal of Mathematics, Vol. IX, Professor Craig has by this method generalized a number of formulae relating to the theory of surfaces. For references to articles on curved spaces will appear in this paper.

\* Sur les propriétés géométriques des courbes gauches dans un espace à  $n$  dimensions - Paris - Gauthier-Villars, 1900.

\*\* Sur les courbes à triple courbure et les surfaces à courbure constante - Paris - Gauthier-Villars, 1900.



I shall confine myself here, for the most part to Riemannian or Euclidean space of four dimensions and shall adopt the nomenclature used by Poincaré in his memoir\* sur les courbes des intégrales doubles and employed by Professor Craig in the paper referred to above. Four dimensional Riemannian space will then be termed hyperspace and a single relation between the coordinates of a point in hyperspace will define a hypersurface two such relations a surface and three a line.

Consider then the hypersurface

$$F(x, y, z, w) = 0$$

of linear element—

\* Acta Math. t. 9. p. 375.





$$ds^2 = dx^2 + dy^2 + dz^2 + dw^2$$

Express the coordinates in terms of three independent parameters. The line element

$$x = x(u, v), \quad y = y(u, v), \quad z = z(u, v), \quad w = w(u, v)$$

which together with the first, second and third partial derivatives, are supposed uniform, finite and continuous through out the region of variation of  $u, v$ .

The linear element then takes the form

$$ds^2 = E_{11} du^2 + E_{22} dv^2 + E_{33} dw^2 + 2E_{12} du dv + 2E_{13} du dw + 2E_{23} dv dw$$

where

$$E_{11} = \sum \left( \frac{\partial x}{\partial u} \right)^2, \quad E_{22} = \sum \left( \frac{\partial x}{\partial v} \right)^2, \quad E_{33} = \sum \left( \frac{\partial x}{\partial w} \right)^2$$

$$E_{12} = \sum \frac{\partial x}{\partial u} \frac{\partial x}{\partial v}, \quad E_{13} = \sum \frac{\partial x}{\partial u} \frac{\partial x}{\partial w}, \quad E_{23} = \sum \frac{\partial x}{\partial v} \frac{\partial x}{\partial w}$$

The study of this ternary differential quadratic form will be the basis of this paper, and after the analogy of the theory of binary quadratic case it is the







$$f = \sum_{i=1}^{u-1} \phi_i du_i du_u$$

where the variables  $u_1, \dots, u_{u-1}$  are functions of  $x_1, \dots, x_n$  and the coefficients  $\phi_i$  are functions of  $u_1, \dots, u_{u-1}$ .

The necessary and sufficient conditions that  $f$  be reducible are\*

$$a = 0$$

$$\sum_i \Delta_{\phi_i} [r^s] = 0$$

where  $a$  is the discriminant of  $f$ ,  $a_{rs}$  the algebraic complement of  $a_{rs}$  in  $a$  and

$$[r^s] = \frac{1}{2} \left( \frac{\partial a_{rs}}{\partial x_r} + \frac{\partial a_{rs}}{\partial x_s} - \frac{\partial a_{rs}}{\partial x_i} \right)$$

Suppose now that  $f$  is irreducible. In this case Schlegel\*\* has shown that it can be deduced from the form

$$f^2 = \sum_{i=1}^h dy_i^2$$

$$\text{where } 0 \leq h \leq \frac{u(u-1)}{2}$$

\* Ricci Loc. Cit. pag. 142.

\*\* Annali di Matematica. (3) vol. 5.



is known as the class\* of the form  $f$ .

Consider the hypersurface

$$x = x(t, u, v), \quad y = y(t, u, v), \quad z = z(t, u, v), \quad w = w(t, u, v)$$

where  $x, y, z, w$  fulfil the conditions prescribed above. If one of the parameters,  $t$  say, is held constant

we get a coordinate surface which I shall call the  $t$  surface.

If two of them,  $u$  and  $v$  say, are held constant we get a coordinate line which I shall call the  $t$  line.

It is to be noted that the  $t$  surface is the surface

$$t = \text{constant}$$

while the  $t$  line is the line

$$u = \text{const.}, \quad v = \text{const.}$$

Jordan\*\* also shows that given the two right bisectors the origin

$$\frac{y'}{a_1} = \dots = \frac{y_n}{a_n}$$

\* Ricci, Loc. Cit., pag. 142.

\*\* Bulletin de la Société Mathématique de France., t. III. pp. 103 —





$$\frac{y}{x} = \dots = \frac{y_n}{x_n}$$

in an  $n$  dimensional space  
the expression

$$\cos^2 \theta = \frac{\sum (a_i a_i')^2}{\sum a_i^2 \sum a_i'^2}$$

is an invariant under orthogonal  
substitution and  $\theta$  is known as  
the angle between the two lines.

$p$  linear equations in the coordinates  
define a linear space of  $n - p$   
dimensions and if we have two  
linear spaces one of  $p$  and the  
other of  $q$  dimensions we can cal-  
culate by means of this formulae  
the angle between any two lines  
one in each space. The expres-  
sion thus obtained is its maxi-  
mum and minimum invariant  
under orthogonal transformation  
and these define the angle be-  
tween the two spaces. The same



The angle between two such lines  
 is the least of the numbers  $\phi g$ ,  
 $u-b$ ,  $u-g$ . Since in hyperspace  
 two hyperplanes make a single  
 angle and if the hyperplanes are  
 written

$$a_1 y_1 + a_2 y_2 + a_3 y_3 + a_4 y_4 + \dots = 0$$

$$b_1 y_1 + b_2 y_2 + b_3 y_3 + b_4 y_4 + \dots = 0$$

the angle is given by

$$\cos^2 \theta = \frac{\sum (a_i b_i)^2}{\sum a_i^2 \sum b_i^2}$$

Write  $\omega_{12}$  = angle between  $t$ -line and  $u$ -line

$$\omega_{13} = \text{ " " } t\text{-line } s\text{-line}$$

$$\omega_{13} = \text{ " " } u\text{-line } s\text{-line}$$

Denoting by  $\cos(\theta_x)$  the cosine of the  
 angle between the  $t$ -line and the  
 $x$ -axis, we have

$$\begin{aligned}
 \cos \omega_{12} &= \sum \cos(\theta_x) \cos(\theta_x) = \frac{1}{\sqrt{\sum \epsilon_{11} \sum \epsilon_{22}}} \sum \frac{\partial x}{\partial t} \frac{\partial x}{\partial u} \\
 &= \frac{\sum \epsilon_{12}}{\sqrt{\sum \epsilon_{11} \sum \epsilon_{22}}}
 \end{aligned}$$



Similarly  $\cos \omega_3 = \frac{E_{13}}{\sqrt{E_{11} E_{33}}}$ ,  $\cos \omega_{23} = \frac{E_{23}}{\sqrt{E_{22} E_{33}}}$

Since the necessary and sufficient conditions for the parametric lines form a triple orthogonal system

$$E_{12} = E_{13} = E_{23} = 0$$

and the line element becomes

$$ds^2 = E_{11} dt^2 + E_{22} du^2 + E_{33} dv^2$$

Since the parametric surfaces, taken in pairs, form six angles the ordinary meaning of triple orthogonal system of surfaces is not immediately extensible to hyperspace but in the case where we have a triple orthogonal system of parametric lines these six angles all become right angles and we have a true triple orthogonal system. This is easily seen as follows.

The tangent plane to the surf.





are at the point  $M(x, y, z, w)$  is given by the  $t$  and  $u$  lines. That is by the following direction cosines

$$P \left( \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial x}{\partial t}, \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial y}{\partial t}, \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial z}{\partial t}, \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial w}{\partial t} \right)$$

$$\left( \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial x}{\partial u}, \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial y}{\partial u}, \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial z}{\partial u}, \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial w}{\partial u} \right)$$

Similarly the tangent plane to the  $u$  surface is given by

$$\Pi \left( \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial x}{\partial t}, \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial y}{\partial t}, \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial z}{\partial t}, \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial w}{\partial t} \right)$$

$$\left( \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial x}{\partial v}, \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial y}{\partial v}, \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial z}{\partial v}, \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial w}{\partial v} \right)$$

Applying Jordan's method the two angles between these two planes  $P$  and  $\Pi$  are given by

$$\begin{vmatrix} \cos \theta & \cos \theta \frac{\epsilon_{11}}{\sqrt{\epsilon_{11} \epsilon_{12}}} & 1 & \frac{\epsilon_{13}}{\sqrt{\epsilon_{11} \epsilon_{33}}} \\ \cos \theta \frac{\epsilon_{11}}{\sqrt{\epsilon_{11} \epsilon_{12}}} & \cos \theta & \frac{\epsilon_{12}}{\sqrt{\epsilon_{11} \epsilon_{22}}} & \frac{\epsilon_{23}}{\sqrt{\epsilon_{11} \epsilon_{33}}} \\ 1 & \frac{\epsilon_{12}}{\sqrt{\epsilon_{11} \epsilon_{12}}} & \cos \theta & \cos \theta \frac{\epsilon_{13}}{\sqrt{\epsilon_{11} \epsilon_{33}}} \\ \frac{\epsilon_{13}}{\sqrt{\epsilon_{11} \epsilon_{33}}} & \frac{\epsilon_{23}}{\sqrt{\epsilon_{22} \epsilon_{33}}} & \cos \theta \frac{\epsilon_{13}}{\sqrt{\epsilon_{11} \epsilon_{33}}} & \cos \theta \end{vmatrix} = 0$$



and it is easily seen that if

$$\varepsilon_{11} = \varepsilon_{22} = \varepsilon_{33} = 0$$

the triangles defined by this equation are right angles and similarly for the other pairs of coordinate surfaces.

Let  $C$  be a line tangent to the hyper-surface making angles  $\vartheta_1, \vartheta_2, \vartheta_3$  with the coordinate axes. Then

$$\cos \vartheta_1 = \frac{1}{\sqrt{g_{11}}} \cos(\alpha) \partial(\alpha) = \frac{1}{\sqrt{g_{11}}} \left( \varepsilon_{11} \frac{dt}{\sqrt{g}} + \varepsilon_{12} \frac{du}{\sqrt{g}} + \varepsilon_{13} \frac{dv}{\sqrt{g}} \right)$$

$$\cos \vartheta_2 = \frac{1}{\sqrt{g_{22}}} \left( \varepsilon_{21} \frac{dt}{\sqrt{g}} + \varepsilon_{22} \frac{du}{\sqrt{g}} + \varepsilon_{23} \frac{dv}{\sqrt{g}} \right), \quad \cos \vartheta_3 = \frac{1}{\sqrt{g_{33}}} \left( \varepsilon_{31} \frac{dt}{\sqrt{g}} + \varepsilon_{32} \frac{du}{\sqrt{g}} + \varepsilon_{33} \frac{dv}{\sqrt{g}} \right)$$

Let  $C'$  be a second line and make displacements by  $\delta$ . The angle between  $C$  and  $C'$  is then

$$\begin{aligned} \cos \vartheta &= \frac{1}{\sqrt{g}} \cos(\alpha) \cos(\alpha') \\ &= \varepsilon_{11} \frac{dt}{\sqrt{g}} \frac{\delta t}{\sqrt{g}} + \varepsilon_{22} \frac{du}{\sqrt{g}} \frac{\delta u}{\sqrt{g}} + \varepsilon_{33} \frac{dv}{\sqrt{g}} \frac{\delta v}{\sqrt{g}} \\ &+ \varepsilon_{12} \left( \frac{dt}{\sqrt{g}} \frac{\delta u}{\sqrt{g}} + \frac{\delta t}{\sqrt{g}} \frac{du}{\sqrt{g}} \right) + \varepsilon_{13} \left( \frac{dt}{\sqrt{g}} \frac{\delta v}{\sqrt{g}} + \frac{\delta t}{\sqrt{g}} \frac{dv}{\sqrt{g}} \right) + \varepsilon_{23} \left( \frac{du}{\sqrt{g}} \frac{\delta v}{\sqrt{g}} + \frac{\delta u}{\sqrt{g}} \frac{dv}{\sqrt{g}} \right) \end{aligned}$$

Since the condition that  $C$  and  $C'$  be orthogonal is that the right hand member of this equation vanish.



## Christoffel Symbols

Writing  $x_1 = t$ ,  $x_2 = u$ ,  $x_3 = v$ ,  $e_{ij} = g_{ij}$  and applying the formula defining the symbols of the first kind

$$[\begin{smallmatrix} i \\ l \end{smallmatrix} \kappa] = \frac{1}{2} \left( \frac{\partial e_{i\ell}}{\partial x_\kappa} + \frac{\partial e_{\kappa\ell}}{\partial x_i} - \frac{\partial e_{i\kappa}}{\partial x_\ell} \right)$$

we have the following eighteen symbols of the first kind of the metric for the case of four variables.

$$[\begin{smallmatrix} 1 \\ 1 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{11}}{\partial t}; [\begin{smallmatrix} 1 \\ 1 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{11}}{\partial u}; [\begin{smallmatrix} 1 \\ 1 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{11}}{\partial v};$$

$$[\begin{smallmatrix} 2 \\ 1 \end{smallmatrix} ] = \frac{\partial g_{12}}{\partial u} - \frac{1}{2} \frac{\partial g_{22}}{\partial t}; [\begin{smallmatrix} 2 \\ 1 \end{smallmatrix} ] = \frac{1}{2} \left( \frac{\partial g_{12}}{\partial v} + \frac{\partial g_{13}}{\partial u} - \frac{\partial g_{23}}{\partial t} \right);$$

$$[\begin{smallmatrix} 3 \\ 1 \end{smallmatrix} ] = \frac{\partial g_{13}}{\partial v} - \frac{1}{2} \frac{\partial g_{33}}{\partial t}; [\begin{smallmatrix} 1 \\ 2 \end{smallmatrix} ] = \frac{\partial g_{12}}{\partial t} - \frac{1}{2} \frac{\partial g_{11}}{\partial u};$$

$$[\begin{smallmatrix} 1 \\ 2 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{22}}{\partial t}; [\begin{smallmatrix} 1 \\ 3 \end{smallmatrix} ] = \frac{1}{2} \left( \frac{\partial g_{12}}{\partial v} + \frac{\partial g_{23}}{\partial t} - \frac{\partial g_{13}}{\partial u} \right); [\begin{smallmatrix} 2 \\ 2 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{22}}{\partial u};$$

$$[\begin{smallmatrix} 2 \\ 2 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{22}}{\partial v}; [\begin{smallmatrix} 3 \\ 3 \end{smallmatrix} ] = \frac{\partial g_{23}}{\partial v} - \frac{1}{2} \frac{\partial g_{33}}{\partial u}; [\begin{smallmatrix} 1 \\ 3 \end{smallmatrix} ] = \frac{\partial g_{13}}{\partial t} - \frac{1}{2} \frac{\partial g_{11}}{\partial v};$$

$$[\begin{smallmatrix} 3 \\ 1 \end{smallmatrix} ] = \frac{1}{2} \left( \frac{\partial g_{13}}{\partial u} + \frac{\partial g_{23}}{\partial t} - \frac{\partial g_{12}}{\partial v} \right); [\begin{smallmatrix} 1 \\ 3 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{33}}{\partial t};$$

$$[\begin{smallmatrix} 2 \\ 3 \end{smallmatrix} ] = \frac{\partial g_{23}}{\partial u} - \frac{1}{2} \frac{\partial g_{22}}{\partial v}; [\begin{smallmatrix} 2 \\ 3 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{33}}{\partial u}; [\begin{smallmatrix} 3 \\ 3 \end{smallmatrix} ] = \frac{1}{2} \frac{\partial g_{33}}{\partial v}.$$



The symbols of the second kind & the indices will be useful later when the hypersurface is referred to a triple or tetragonal system and I shall calculate them for test case. Taking

$$\mathcal{E}_{12} = \mathcal{E}_{13} = \mathcal{E}_{23} = 0$$

in the symbols of the first kind these are given at once, and yet those of the second kind it is sufficient to observe that

$\tau_{ii} = \frac{1}{\mathcal{E}_{ii}}$ ,  $\tau_{22} = \frac{1}{\mathcal{E}_{22}}$ ,  $\tau_{33} = \frac{1}{\mathcal{E}_{33}}$ ,  $\tau_{ii} = 0$  ( $i \neq l$ ).  
 where  $\tau_{ij}$  is the algebraic complement of  $\mathcal{E}_{ij}$  in  $\Delta$  divided by  $\Delta$  itself.

Computing now the symbols of the second kind from the formula

$$\{i^r\} = \sum \tau_{rl} [i^l]$$

we have

$$\{1^1\} = \frac{1}{\mathcal{E}_{11}} \frac{\sqrt{\mathcal{E}_{11}}}{\sqrt{\Delta}}, \quad \{2^2\} = \frac{1}{\mathcal{E}_{22}} \frac{\sqrt{\mathcal{E}_{22}}}{\sqrt{\Delta}}, \quad \{3^3\} = \frac{1}{\mathcal{E}_{33}} \frac{\sqrt{\mathcal{E}_{33}}}{\sqrt{\Delta}},$$

$$\{1^2\} = -\frac{\mathcal{E}_{12}}{\mathcal{E}_{11}} \frac{\sqrt{\mathcal{E}_{12}}}{\sqrt{\Delta}}, \quad \{2^1\} = 0, \quad \{3^1\} = -\frac{\mathcal{E}_{13}}{\mathcal{E}_{11}} \frac{\sqrt{\mathcal{E}_{33}}}{\sqrt{\Delta}},$$





$$\{1\} = -\frac{\sqrt{\epsilon_{11}}}{\epsilon_{11}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial u}, \quad \{1, 2\} = \frac{1}{\sqrt{\epsilon_{12}}} \frac{\partial \sqrt{\epsilon_{12}}}{\partial \sigma}, \quad \{1, 3\} = 0;$$

$$\{2, 2\} = \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial u}, \quad \{2, 3\} = \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial \sigma}, \quad \{3, 3\} = -\frac{\sqrt{\epsilon_{33}}}{\epsilon_{22}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial u}$$

$$\{1, 1\} = -\frac{\sqrt{\epsilon_{11}}}{\epsilon_{33}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial \sigma}, \quad \{1, 2\} = 0, \quad \{1, 3\} = \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial u}$$

$$\{2, 2\} = -\frac{\sqrt{\epsilon_{22}}}{\epsilon_{33}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial \sigma}, \quad \{2, 3\} = \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial u}, \quad \{3, 3\} = \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial \sigma}$$

Quote now the discriminant of the linear element by  $u$  when expressed in the  $a_{ij}$  and by  $\sigma$  when expressed in the  $\epsilon_{ij}$  and calculate the differentiated parameters  $1\phi$ ,  $1_2\phi$ ,  $1(\phi\psi)$  from the formulae

$$1\phi = \frac{3}{2s} \epsilon_{rs} \frac{\partial \phi}{\partial x_r} \frac{\partial \phi}{\partial x_s}$$

$$1_2\phi = \frac{3}{1a} \frac{\partial}{\partial x_s} \frac{3}{2r} \epsilon_{rs} \frac{\partial \phi}{\partial x_r}$$

$$1(\phi\psi) = \frac{3}{2rs} \epsilon_{rs} \frac{\partial \phi}{\partial x_r} \frac{\partial \psi}{\partial x_s}$$

The necessary and sufficient condition that the two surfaces  $\phi = \text{const.}$   $\psi = \text{const.}$



orthogonal is here expressed by\*

$$1(\phi\psi) = 0$$

## II

The tangent hyperplane is determined by any three lines lying in it and in particular by the tangents to the three coordinate lines. Defining then the normal to the hypersurface at any point  $M$  as the line perpendicular to the hypertangent plane at  $M$ , we have & express that the line whose direction cosines are  $X, Y, Z, W$  is orthogonal to the three coordinate lines.

We have then

$$\sum X \frac{\partial x}{\partial t} = 0, \quad \sum X \frac{\partial x}{\partial u} = 0, \quad \sum X \frac{\partial x}{\partial v} = 0$$

These give for the direction cosines of the normal

$$X = \frac{1}{\sqrt{A}} \frac{\partial(yzw)}{\partial(tuv)}, \quad Y = \frac{1}{\sqrt{A}} \frac{\partial(zwx)}{\partial(tuv)}, \quad Z = \frac{1}{\sqrt{A}} \frac{\partial(wxy)}{\partial(tuv)}, \quad W = \frac{1}{\sqrt{A}} \frac{\partial(xyz)}{\partial(tuv)} \quad (1)$$

\* Padoa - with three like coordinate-curvatures - Rend. della R. Acc. dei Lincei. 1907, p. 115.



Introduce now the second differential quadratic form

$$\begin{aligned} \varphi &= -\sum dX dx \\ &= J_1 dt^2 + J_2 du^2 + J_3 dv^2 + 2J_2 dt du + 2J_3 dt dv + 2J_{23} du dv \end{aligned} \quad (2)$$

We at once find

$$\left. \begin{aligned} J_1 &= \sum X \frac{\partial^2 x}{\partial t^2} = -\sum \frac{\partial X}{\partial t} \frac{\partial x}{\partial t} \\ J_2 &= \sum X \frac{\partial^2 x}{\partial t \partial u} = -\sum \frac{\partial X}{\partial t} \frac{\partial x}{\partial u} \\ J_3 &= \sum X \frac{\partial^2 x}{\partial t \partial v} = -\sum \frac{\partial X}{\partial t} \frac{\partial x}{\partial v} \\ J_2 &= \sum X \frac{\partial^2 x}{\partial t \partial u} = -\sum \frac{\partial X}{\partial t} \frac{\partial x}{\partial u} = -\sum \frac{\partial X}{\partial u} \frac{\partial x}{\partial t} \\ J_3 &= \sum X \frac{\partial^2 x}{\partial t \partial v} = -\sum \frac{\partial X}{\partial t} \frac{\partial x}{\partial v} = -\sum \frac{\partial X}{\partial v} \frac{\partial x}{\partial t} \\ J_{23} &= \sum X \frac{\partial^2 x}{\partial u \partial v} = -\sum \frac{\partial X}{\partial u} \frac{\partial x}{\partial v} = -\sum \frac{\partial X}{\partial v} \frac{\partial x}{\partial u} \end{aligned} \right\} \quad (3)$$

By means of the expressions (\*) these can be written

$$J_1 = \frac{1}{\Delta} \begin{vmatrix} \frac{\partial y}{\partial t} & \frac{\partial z}{\partial t} & \frac{\partial w}{\partial t} \\ \frac{\partial x}{\partial u} & \frac{\partial y}{\partial u} & \frac{\partial z}{\partial u} \\ \frac{\partial x}{\partial v} & \frac{\partial y}{\partial v} & \frac{\partial z}{\partial v} \end{vmatrix}$$

with similar expressions for the remaining five.



of four  $x, y, z, w$  or any four functions whatever of  $t, u, v, w, x, y, z, w$  can be so determined that the following four equations are satisfied

$$\left. \begin{aligned} A &= a \frac{dx}{dt} + b \frac{dx}{du} + c \frac{dx}{dv} + d \frac{dx}{dw} \\ B &= a \frac{dy}{dt} + b \frac{dy}{du} + c \frac{dy}{dv} + d \frac{dy}{dw} \\ C &= a \frac{dz}{dt} + b \frac{dz}{du} + c \frac{dz}{dv} + d \frac{dz}{dw} \\ D &= a \frac{dw}{dt} + b \frac{dw}{du} + c \frac{dw}{dv} + d \frac{dw}{dw} \end{aligned} \right\} \quad (a)$$

since the determinant of the system equals 14 and is therefore not zero.

By exactly the same reasoning employed in the case of two parameters we arrive at the following system of equations satisfied by the four coordinates  $x, y, z, w$ .

$$\left. \begin{aligned} \frac{\partial A}{\partial t} &= \left\{ \frac{\partial}{\partial t} \right\} \frac{\partial \theta}{\partial t} + \left\{ \frac{\partial}{\partial u} \right\} \frac{\partial \theta}{\partial u} + \left\{ \frac{\partial}{\partial v} \right\} \frac{\partial \theta}{\partial v} + \left\{ \frac{\partial}{\partial w} \right\} \frac{\partial \theta}{\partial w} \\ \frac{\partial A}{\partial u} &= \left\{ \frac{\partial}{\partial t} \right\} \frac{\partial \theta}{\partial u} + \left\{ \frac{\partial}{\partial u} \right\} \frac{\partial \theta}{\partial u} + \left\{ \frac{\partial}{\partial v} \right\} \frac{\partial \theta}{\partial v} + \left\{ \frac{\partial}{\partial w} \right\} \frac{\partial \theta}{\partial w} \end{aligned} \right\}$$

+ Bianchi - Lezioni di Geometria  
Differenziale - 6.57





$$\left. \begin{aligned} \frac{\partial^2 \theta}{\partial t^2} &= \left\{ \begin{matrix} 1 \\ 1 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial t^2} + \left\{ \begin{matrix} 2 \\ 1 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial u \partial t} + \left\{ \begin{matrix} 3 \\ 1 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial v \partial t} + \Delta_{13} X \\ \frac{\partial^2 \theta}{\partial u^2} &= \left\{ \begin{matrix} 2 \\ 1 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial t^2} + \left\{ \begin{matrix} 2 \\ 2 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial u \partial t} + \left\{ \begin{matrix} 3 \\ 2 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial v \partial t} + \Delta_{22} X \\ \frac{\partial^2 \theta}{\partial u \partial v} &= \left\{ \begin{matrix} 3 \\ 1 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial t^2} + \left\{ \begin{matrix} 3 \\ 2 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial u \partial t} + \left\{ \begin{matrix} 3 \\ 3 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial v \partial t} + \Delta_{33} X \\ \frac{\partial^2 \theta}{\partial v^2} &= \left\{ \begin{matrix} 3 \\ 1 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial t^2} + \left\{ \begin{matrix} 3 \\ 2 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial u \partial t} + \left\{ \begin{matrix} 3 \\ 3 \end{matrix} \right\} \frac{\partial^2 \theta}{\partial v \partial t} + \Delta_{33} X \end{aligned} \right\} \quad (I)$$

and the three equations satisfied by the direction cosines of the normal

$$\left. \begin{aligned} \frac{\partial \Theta}{\partial t} &= -\frac{1}{\Delta} |\Delta'_{11}| \frac{\partial \theta}{\partial t} - \frac{1}{\Delta} |\Delta'_{12}| \frac{\partial \theta}{\partial u} - \frac{1}{\Delta} |\Delta'_{13}| \frac{\partial \theta}{\partial v} \\ \frac{\partial \Theta}{\partial u} &= -\frac{1}{\Delta} |\Delta'_{21}| \frac{\partial \theta}{\partial t} - \frac{1}{\Delta} |\Delta'_{22}| \frac{\partial \theta}{\partial u} - \frac{1}{\Delta} |\Delta'_{23}| \frac{\partial \theta}{\partial v} \\ \frac{\partial \Theta}{\partial v} &= -\frac{1}{\Delta} |\Delta'_{31}| \frac{\partial \theta}{\partial t} - \frac{1}{\Delta} |\Delta'_{32}| \frac{\partial \theta}{\partial u} - \frac{1}{\Delta} |\Delta'_{33}| \frac{\partial \theta}{\partial v} \end{aligned} \right\} \quad (II)$$

where

$$|\Delta'_{i1}| \equiv \begin{vmatrix} \Delta_{i1} & \Delta_{i2} & \Delta_{i3} \\ \Delta_{21} & \Delta_{22} & \Delta_{23} \\ \Delta_{31} & \Delta_{32} & \Delta_{33} \end{vmatrix} \quad (i=1, 2, 3)$$

A column of  $\Delta_{i1}$  standing in the first, second or third place according to the upper index hence in the first place in this case.



The relations connecting the coefficients of the first and second forms can now be found by writing the conditions that the two systems of partial differential equations (I) and (II) be integrable.

Putting as usual  $x_1, x_2, x_3$  for the coordinates  $u, v$  the conditions of integrability are for (I)

$$\frac{\partial}{\partial x_t} \left( \frac{\partial^2 \theta}{\partial x_r \partial x_s} \right) = \frac{\partial}{\partial x_r} \left( \frac{\partial^2 \theta}{\partial x_s \partial x_t} \right)$$

$$t, r, s = 1, 2, 3, \quad t \neq r, \quad t \neq s.$$

These give

$$\sum_i \left[ \{rs\}_i \frac{\partial^2 \theta}{\partial x_i \partial x_t} + \frac{\partial}{\partial x_t} \{i\}_s \frac{\partial \theta}{\partial x_i} + \odot \frac{\partial \partial_{rs}}{\partial x_t} + \partial_{rs} \frac{\partial \odot}{\partial x_t} \right] =$$

$$\sum_i \left[ \{st\}_i \frac{\partial^2 \theta}{\partial x_i \partial x_r} + \frac{\partial}{\partial x_r} \{i\}_s \frac{\partial \theta}{\partial x_i} + \odot \frac{\partial \partial_{st}}{\partial x_r} + \partial_{st} \frac{\partial \odot}{\partial x_r} \right]$$

Substituting for  $\frac{\partial^2 \theta}{\partial x_r \partial x_s}$  and  $\frac{\partial \odot}{\partial x_t}$  from equations (I) and (II) we have



$$\sum_{ij} \left[ \left\{ \begin{matrix} rs \\ ij \end{matrix} \right\} \left\{ \begin{matrix} it \\ js \end{matrix} \right\} - \left\{ \begin{matrix} st \\ ir \end{matrix} \right\} \left\{ \begin{matrix} jr \\ is \end{matrix} \right\} + \frac{\partial}{\partial x_s} \left\{ \begin{matrix} rs \\ js \end{matrix} \right\} - \frac{\partial}{\partial x_r} \left\{ \begin{matrix} st \\ is \end{matrix} \right\} - \right. \\ \left. \frac{\partial s}{\partial t} \left| \mathcal{D}_{ij}^i \right| + \frac{\partial t}{\partial s} \left| \mathcal{D}_{ij}^j \right| \right) \frac{\partial \theta}{\partial x_j} +$$

$$\left( \frac{\partial \mathcal{D}_{rs}}{\partial x_s} - \frac{\partial \mathcal{D}_{st}}{\partial x_r} + \mathcal{D}_{st} \left\{ \begin{matrix} rs \\ js \end{matrix} \right\} - \mathcal{D}_{rs} \left\{ \begin{matrix} st \\ is \end{matrix} \right\} \right) \mathcal{C} = 0 \quad (b)$$

substituting  $x, y, z, w$  successively for  $\theta$  we have a system of four equations whose determinants

$$\begin{vmatrix} \frac{\partial x}{\partial t} & \frac{\partial x}{\partial u} & \frac{\partial x}{\partial v} & X \\ \frac{\partial y}{\partial t} & \frac{\partial y}{\partial u} & \frac{\partial y}{\partial v} & Y \\ \frac{\partial z}{\partial t} & \frac{\partial z}{\partial u} & \frac{\partial z}{\partial v} & Z \\ \frac{\partial w}{\partial t} & \frac{\partial w}{\partial u} & \frac{\partial w}{\partial v} & W \end{vmatrix} = \sqrt{A}$$

do not vanish, and hence the coefficients must separately vanish. Introducing here the <sup>Poisson</sup> ~~Christoffel~~ symbols of four indices

$$\left\{ \begin{matrix} s \\ j \end{matrix} \right\} \left\{ \begin{matrix} t \\ i \end{matrix} \right\} = \frac{\partial}{\partial x_s} \left\{ \begin{matrix} rs \\ js \end{matrix} \right\} - \frac{\partial}{\partial x_r} \left\{ \begin{matrix} st \\ is \end{matrix} \right\} + \sum_{ij} \left( \left\{ \begin{matrix} rs \\ ij \end{matrix} \right\} \left\{ \begin{matrix} it \\ js \end{matrix} \right\} - \left\{ \begin{matrix} st \\ ir \end{matrix} \right\} \left\{ \begin{matrix} jr \\ is \end{matrix} \right\} \right)$$

we have the equations,



$$\{sj \ r t\} - \frac{\partial_{rs}}{\Delta} |\partial_{rt}^j| + \frac{\partial_{rt}}{\Delta} |\partial_{rs}^j| = 0 \quad (III)$$

$$\frac{\partial \partial_{rs}}{\partial x_t} - \frac{\partial \partial_{st}}{\partial x_r} + \Delta (\partial_{rt} \{j \ r s\} - \partial_{rs} \{j \ r t\}) = 0 \quad (IV)$$

$t, r, s = 1, 2, 3; t \neq r; t \neq s; j = 1, 2, 3.$

From (III) we derive the twenty four equations

$$\left. \begin{aligned} \{1j \ 12\} - \frac{\partial_{11}}{\Delta} |\partial_{12}^j| + \frac{\partial_{12}}{\Delta} |\partial_{11}^j| &= 0 \\ \{1j \ 13\} - \frac{\partial_{11}}{\Delta} |\partial_{13}^j| + \frac{\partial_{13}}{\Delta} |\partial_{11}^j| &= 0 \\ \{2j \ 11\} - \frac{\partial_{22}}{\Delta} |\partial_{11}^j| + \frac{\partial_{12}}{\Delta} |\partial_{22}^j| &= 0 \\ \{1j \ 23\} - \frac{\partial_{22}}{\Delta} |\partial_{13}^j| + \frac{\partial_{23}}{\Delta} |\partial_{12}^j| &= 0 \\ \{3j \ 11\} - \frac{\partial_{33}}{\Delta} |\partial_{11}^j| + \frac{\partial_{13}}{\Delta} |\partial_{33}^j| &= 0 \\ \{3j \ 32\} - \frac{\partial_{33}}{\Delta} |\partial_{12}^j| + \frac{\partial_{23}}{\Delta} |\partial_{33}^j| &= 0 \\ \{3j \ 21\} - \frac{\partial_{23}}{\Delta} |\partial_{11}^j| + \frac{\partial_{13}}{\Delta} |\partial_{22}^j| &= 0 \\ \{2j \ 31\} - \frac{\partial_{23}}{\Delta} |\partial_{11}^j| + \frac{\partial_{12}}{\Delta} |\partial_{33}^j| &= 0 \end{aligned} \right\} \quad (III)$$

which are a generalization of the four\* equations of the theory of surfaces giving the curvature of the first fundamental form.

\* Bianchi - Lezioni p. 90.





From II I get the eight equations

$$\frac{\partial \lambda_{11}}{\partial u} - \frac{\partial \lambda_{12}}{\partial t} + \left\{ \begin{matrix} 11 \\ 22 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 12 \\ 22 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 13 \\ 22 \end{matrix} \right\} \lambda_{13} + \left\{ \begin{matrix} 21 \\ 22 \end{matrix} \right\} \lambda_{21} + \left\{ \begin{matrix} 22 \\ 22 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 23 \\ 22 \end{matrix} \right\} \lambda_{23} = 0$$

$$\frac{\partial \lambda_{11}}{\partial v} - \frac{\partial \lambda_{13}}{\partial t} + \left\{ \begin{matrix} 11 \\ 32 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 12 \\ 32 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 13 \\ 32 \end{matrix} \right\} \lambda_{13} + \left\{ \begin{matrix} 21 \\ 32 \end{matrix} \right\} \lambda_{21} + \left\{ \begin{matrix} 22 \\ 32 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 23 \\ 32 \end{matrix} \right\} \lambda_{23} = 0$$

$$\frac{\partial \lambda_{22}}{\partial t} - \frac{\partial \lambda_{12}}{\partial u} + \left\{ \begin{matrix} 22 \\ 11 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 22 \\ 22 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 22 \\ 32 \end{matrix} \right\} \lambda_{23} + \left\{ \begin{matrix} 11 \\ 22 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 12 \\ 22 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 13 \\ 22 \end{matrix} \right\} \lambda_{13} = 0$$

$$\frac{\partial \lambda_{22}}{\partial v} - \frac{\partial \lambda_{23}}{\partial u} + \left\{ \begin{matrix} 22 \\ 11 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 22 \\ 22 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 22 \\ 32 \end{matrix} \right\} \lambda_{23} + \left\{ \begin{matrix} 21 \\ 11 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 22 \\ 11 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 23 \\ 11 \end{matrix} \right\} \lambda_{13} = 0$$

$$\left. \begin{aligned} \frac{\partial \lambda_{33}}{\partial t} - \frac{\partial \lambda_{13}}{\partial v} + \left\{ \begin{matrix} 33 \\ 11 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 33 \\ 22 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 33 \\ 32 \end{matrix} \right\} \lambda_{23} + \left\{ \begin{matrix} 11 \\ 33 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 12 \\ 33 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 13 \\ 33 \end{matrix} \right\} \lambda_{13} = 0 \\ \frac{\partial \lambda_{33}}{\partial u} - \frac{\partial \lambda_{23}}{\partial v} + \left\{ \begin{matrix} 33 \\ 21 \end{matrix} \right\} \lambda_{21} + \left\{ \begin{matrix} 33 \\ 22 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 33 \\ 23 \end{matrix} \right\} \lambda_{23} + \left\{ \begin{matrix} 21 \\ 33 \end{matrix} \right\} \lambda_{21} + \left\{ \begin{matrix} 22 \\ 33 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 23 \\ 33 \end{matrix} \right\} \lambda_{23} = 0 \end{aligned} \right\} \text{(II)}$$

$$\frac{\partial \lambda_{33}}{\partial u} - \frac{\partial \lambda_{23}}{\partial v} + \left\{ \begin{matrix} 33 \\ 11 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 33 \\ 22 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 33 \\ 32 \end{matrix} \right\} \lambda_{23} + \left\{ \begin{matrix} 21 \\ 33 \end{matrix} \right\} \lambda_{21} + \left\{ \begin{matrix} 22 \\ 33 \end{matrix} \right\} \lambda_{22} + \left\{ \begin{matrix} 23 \\ 33 \end{matrix} \right\} \lambda_{23} = 0$$

$$\frac{\partial \lambda_{23}}{\partial t} - \frac{\partial \lambda_{13}}{\partial u} + \left\{ \begin{matrix} 23 \\ 11 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 23 \\ 22 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 23 \\ 32 \end{matrix} \right\} \lambda_{13} + \left\{ \begin{matrix} 11 \\ 23 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 12 \\ 23 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 13 \\ 23 \end{matrix} \right\} \lambda_{13} = 0$$

$$\frac{\partial \lambda_{23}}{\partial v} - \frac{\partial \lambda_{12}}{\partial v} + \left\{ \begin{matrix} 23 \\ 11 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 23 \\ 22 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 23 \\ 32 \end{matrix} \right\} \lambda_{13} + \left\{ \begin{matrix} 11 \\ 23 \end{matrix} \right\} \lambda_{11} + \left\{ \begin{matrix} 12 \\ 23 \end{matrix} \right\} \lambda_{12} + \left\{ \begin{matrix} 13 \\ 23 \end{matrix} \right\} \lambda_{13} = 0$$

which are a generalization of the Bianchi-Eddington equations.

Since, given the differential quadratic form

$$I = E_{11} dt^2 + E_{22} du^2 + E_{33} dv^2 + \dots$$

which is definite and whose coefficients satisfy (III) and (II) there



is a hypersurface admitting this as a fixed element.

Ricci\* has deduced the necessary and sufficient conditions that a differential quadratic form in  $n$  variables be irreducible and of the first class, and it is interesting to note that these are exactly the conditions (III) and (IV) for the case of three variables.

Introduce the Christoffel symbols of four indices of the first kind by

$$(rkih) = \sum \alpha_{rk} \{rvih\}.$$

These enjoy the following properties\*\*

$$(rkih) = -(krih)$$

$$(rkih) = -(rkhi)$$

$$(rkih) = (ihrk)$$

$$(rkih) + (rkhi) + (khi r) = 0$$

Taking account of these relations there are six symbols of the first kind  
\* loc. cit.

\*\* Bianchi - Lezioni, p. 50.



not identically zero:

(212), (1213), (1223), (1313), (1323), (2323).

Introduce the symbol of two indices

$$(lq)_{ra} = \begin{vmatrix} \frac{\partial^2 \varphi_1}{\partial x_1 \partial x_1} & \frac{\partial^2 \varphi_1}{\partial x_1 \partial x_2} & \frac{\partial^2 \varphi_1}{\partial x_1 \partial x_3} & \frac{\partial^2 \varphi_1}{\partial x_1 \partial x_4} \\ \frac{\partial^2 \varphi_2}{\partial x_2 \partial x_1} & \frac{\partial^2 \varphi_2}{\partial x_2 \partial x_2} & \frac{\partial^2 \varphi_2}{\partial x_2 \partial x_3} & \frac{\partial^2 \varphi_2}{\partial x_2 \partial x_4} \\ \frac{\partial^2 \varphi_3}{\partial x_3 \partial x_1} & \frac{\partial^2 \varphi_3}{\partial x_3 \partial x_2} & \frac{\partial^2 \varphi_3}{\partial x_3 \partial x_3} & \frac{\partial^2 \varphi_3}{\partial x_3 \partial x_4} \\ \frac{\partial^2 \varphi_4}{\partial x_4 \partial x_1} & \frac{\partial^2 \varphi_4}{\partial x_4 \partial x_2} & \frac{\partial^2 \varphi_4}{\partial x_4 \partial x_3} & \frac{\partial^2 \varphi_4}{\partial x_4 \partial x_4} \end{vmatrix}$$

and the symbol of three indices

$$(lmr) = \frac{\partial(lp)}{\partial x_m} - \frac{\partial(lm)}{\partial x_p} + \sum_{rs} \epsilon_{rsg} [(rm)[s] - (rp)[s]]$$

Then the necessary and sufficient conditions that the differential quadratic form

$$f = \sum_{rs} a_{rs} dx_r dx_s$$

is irreducible and of first class are

$$a \neq 0; (lmr) = (lp)_{rm} - (lp)_{mr}, (lmr) = 0$$

3

To identify these with equations (I) and (II) we now simply write

$$x = x_1, y = x_2, z = x_3, u = x_4, v = x_5, w = x_6, t = x_7$$

and we see at once from (3\*) p. 17 that  $(g) = d_{ij}$  coefficient of the second form since the symbols  $(\text{sup } g)$  are the minors of the second order in the discriminant of the second form

$$D = \begin{vmatrix} d_{11} & d_{12} & d_{13} \\ d_{21} & d_{22} & d_{23} \\ d_{31} & d_{32} & d_{33} \end{vmatrix}$$

Solving the system of equations

$$(rkih) = \sum a_{rk} \{rvkh\}$$

we have

$$\{rvkh\} = \sum A_{rk} (rkih)$$

and substituting

$$(rkih) = (r_i)(k_h) - (r_h)(k_i)$$

we have at once

$$\{rvkh\} - \sum A_{rk} (r_i)(k_h) + \sum A_{rk} (r_h)(k_i) = 0$$

or





$$\{r, i\} - \frac{\partial r}{\partial x_i} |D_{\mu}^r| + \frac{\partial r}{\partial x} |D_{\mu}^r| = 0 \quad (c)$$

Substituting  $D_p = (p)$  in  
 $(p, p) = 0$

we get

$$\frac{\partial D_p}{\partial x_{\mu}} - \frac{\partial D_{\mu}}{\partial x_p} + \sum_r (D_{\mu} \{p\}^r - D_p \{r\}^{\mu}) = 0 \quad (d)$$

These equations (c) and (d) are exactly equations (III) and (IV) and we may say:

In order that a differential  
algebraic expression in the  
variables represent the linear ele-  
ment of a hypersurface it must  
be irreducible and of the first class.

Equations (III) and (IV) can be written in a form which will be useful when the existence of a hypersurface corresponding to the two fundamental forms is considered.

From (III) we have at once



$$\frac{1}{\sqrt{a}} \frac{\partial \rho_p}{\partial x_m} - \frac{1}{\sqrt{a}} \frac{\partial \rho_m}{\partial x_p} + \frac{1}{\sqrt{a}} [\rho_m \{l^p\} - \rho_p \{l^m\}] + \frac{1}{\sqrt{a}} \Sigma' [\rho_m \{r\}^p - \rho_p \{r\}^m] = 0 \quad (a)$$

where in  $\Sigma'$   $r=l$  is excluded.

$$\text{Now } \frac{\partial \log \sqrt{a}}{\partial x_i} = \Sigma' \sin [i^l]$$

$$\text{Hence } \frac{\partial (\frac{1}{\sqrt{a}})}{\partial x_i} = \Sigma' \frac{1}{\sqrt{a}} \{i^l\}$$

and we have

$$\rho_p \frac{\partial (\frac{1}{\sqrt{a}})}{\partial x_m} - \rho_m \frac{\partial (\frac{1}{\sqrt{a}})}{\partial x_p} + \frac{1}{\sqrt{a}} (\rho_p \{l^m\} - \rho_m \{l^p\}) + \frac{1}{\sqrt{a}} \Sigma' (\rho_p \{r\}^m - \rho_m \{r\}^p) = 0 \quad (b)$$

Adding (a) and (b)

$$\frac{\partial (\frac{\rho_p}{\sqrt{a}})}{\partial x_m} - \frac{\partial (\frac{\rho_m}{\sqrt{a}})}{\partial x_p} + \frac{1}{\sqrt{a}} \Sigma' (\rho_m \{r\}^p - \rho_p \{r\}^m) + (\rho_p \{l^m\} - \rho_m \{l^p\}) = 0 \quad (I)$$

These hold whatever be the co-ordinate lines. Suppose the hypersurfaces are referred to a triple orthogonal system, which at the same time reduces the manifold to a sum of squares.



It will appear later that the system formed by the lines of curvature in this case three out of the independent symbols ( $\Gamma_{ij}^k$ ) are different from zero, namely

$$\{\Gamma_{ij}^i\} = \Gamma_{ij}^i \quad (i, j = 1, 2, 3, i \neq j)$$

The remaining three, not identically zero are

$$\{\Gamma_{ij}^k\} = 0 \quad (i, j, k = 1, 2, 3, i \neq j \neq k)$$

$$\begin{aligned} \{\Gamma_{ij}^k\} &= \sum_r \Gamma_{rj}^k \Gamma_{ir}^i \\ &= \Gamma_{ij}^k \end{aligned}$$

It is at once

$$\{\Gamma_{ij}^i\} = \frac{\partial \Gamma_{ij}^i}{\partial x^j}, \quad \{\Gamma_{ji}^j\} = \frac{\partial \Gamma_{ji}^j}{\partial x^i}$$

$$\{\Gamma_{ij}^k\} = 0 \quad (i \neq j \neq k)$$

Equations (III) then become

$$\{\Gamma_{ij}^i\} = \frac{\partial \Gamma_{ij}^i}{\partial x^j}, \quad \{\Gamma_{ij}^k\} = 0 \quad (i \neq j \neq k) \quad (III^*)$$

and equations (II) take the form



$$\frac{\partial \left( \frac{\lambda_{uu}}{r\bar{a}} \right)}{\partial \lambda_{uu}} + \frac{\lambda_{uu}}{r\bar{a}} \{ \lambda \} = 0$$

(II\*)

Calculating these we have

$$\frac{\partial}{\partial t} \left( \frac{1}{\bar{\epsilon}_{11}} \frac{\partial \bar{\epsilon}_{11}}{\partial t} \right) + \frac{\partial}{\partial u} \left( \frac{1}{\bar{\epsilon}_{11}} \frac{\partial \bar{\epsilon}_{11}}{\partial u} \right) + \frac{1}{\bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{11}}{\partial v} \frac{\partial \bar{\epsilon}_{11}}{\partial v} + \frac{\lambda_{11} \lambda_{11}}{\bar{\epsilon}_{11} \bar{\epsilon}_{11}} = 0$$

$$\frac{\partial}{\partial t} \left( \frac{1}{\bar{\epsilon}_{11}} \frac{\partial \bar{\epsilon}_{33}}{\partial t} \right) + \frac{\partial}{\partial v} \left( \frac{1}{\bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{11}}{\partial v} \right) - \frac{1}{\bar{\epsilon}_{22}} \frac{\partial \bar{\epsilon}_{11}}{\partial u} \frac{\partial \bar{\epsilon}_{33}}{\partial u} + \frac{\lambda_{11} \lambda_{33}}{\bar{\epsilon}_{11} \bar{\epsilon}_{33}} = 0$$

$$\frac{\partial}{\partial u} \left( \frac{1}{\bar{\epsilon}_{22}} \frac{\partial \bar{\epsilon}_{33}}{\partial u} \right) + \frac{\partial}{\partial v} \left( \frac{1}{\bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{22}}{\partial v} \right) + \frac{1}{\bar{\epsilon}_{11}} \frac{\partial \bar{\epsilon}_{22}}{\partial t} \frac{\partial \bar{\epsilon}_{33}}{\partial t} + \frac{\lambda_{22} \lambda_{33}}{\bar{\epsilon}_{22} \bar{\epsilon}_{33}} = 0$$

$$\frac{\partial^2 \bar{\epsilon}_{33}}{\partial t \partial u} - \frac{1}{\bar{\epsilon}_{11}} \frac{\partial \bar{\epsilon}_{11}}{\partial u} \frac{\partial \bar{\epsilon}_{33}}{\partial t} - \frac{1}{\bar{\epsilon}_{22}} \frac{\partial \bar{\epsilon}_{22}}{\partial t} \frac{\partial \bar{\epsilon}_{33}}{\partial u} = 0$$

$$\frac{\partial^2 \bar{\epsilon}_{11}}{\partial t \partial v} - \frac{1}{\bar{\epsilon}_{11}} \frac{\partial \bar{\epsilon}_{11}}{\partial v} \frac{\partial \bar{\epsilon}_{11}}{\partial t} - \frac{1}{\bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{22}}{\partial v} \frac{\partial \bar{\epsilon}_{33}}{\partial t} = 0$$

$$\frac{\partial^2 \bar{\epsilon}_{11}}{\partial u \partial v} - \frac{1}{\bar{\epsilon}_{22}} \frac{\partial \bar{\epsilon}_{11}}{\partial u} \frac{\partial \bar{\epsilon}_{22}}{\partial v} - \frac{1}{\bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{22}}{\partial v} \frac{\partial \bar{\epsilon}_{33}}{\partial u} = 0$$

The last three of these may also be written

$$\frac{\partial}{\partial u} \left( \frac{1}{\bar{\epsilon}_{11}} \frac{\partial \bar{\epsilon}_{33}}{\partial t} \right) - \frac{1}{\bar{\epsilon}_{11} \bar{\epsilon}_{22}} \frac{\partial \bar{\epsilon}_{11}}{\partial t} \frac{\partial \bar{\epsilon}_{33}}{\partial u} = 0$$

$$\frac{\partial}{\partial t} \left( \frac{1}{\bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{22}}{\partial v} \right) - \frac{1}{\bar{\epsilon}_{11} \bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{11}}{\partial v} \frac{\partial \bar{\epsilon}_{22}}{\partial t} = 0$$

$$\frac{\partial}{\partial u} \left( \frac{1}{\bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{11}}{\partial v} \right) - \frac{1}{\bar{\epsilon}_{11} \bar{\epsilon}_{33}} \frac{\partial \bar{\epsilon}_{11}}{\partial u} \frac{\partial \bar{\epsilon}_{22}}{\partial v} = 0$$

(III\*)





From (IX\*) we have

$$\left. \begin{aligned} \frac{\partial}{\partial u} \left( \frac{\nu_{11}}{\sqrt{\Sigma_{11}}} \right) - \frac{\nu_{12}}{\Sigma_{12}} \frac{\partial \sqrt{\Sigma_{11}}}{\partial u} &= 0 \\ \frac{\partial}{\partial v} \left( \frac{\nu_{11}}{\sqrt{\Sigma_{11}}} \right) - \frac{\nu_{33}}{\Sigma_{33}} \frac{\partial \sqrt{\Sigma_{11}}}{\partial v} &= 0 \\ \frac{\partial}{\partial t} \left( \frac{\nu_{22}}{\sqrt{\Sigma_{22}}} \right) - \frac{\nu_{11}}{\Sigma_{11}} \frac{\partial \sqrt{\Sigma_{22}}}{\partial t} &= 0 \\ \frac{\partial}{\partial r} \left( \frac{\nu_{22}}{\sqrt{\Sigma_{22}}} \right) - \frac{\nu_{33}}{\Sigma_{33}} \frac{\partial \sqrt{\Sigma_{22}}}{\partial r} &= 0 \\ \frac{\partial}{\partial t} \left( \frac{\nu_{33}}{\sqrt{\Sigma_{33}}} \right) - \frac{\nu_{11}}{\Sigma_{11}} \frac{\partial \sqrt{\Sigma_{33}}}{\partial t} &= 0 \\ \frac{\partial}{\partial u} \left( \frac{\nu_{33}}{\sqrt{\Sigma_{33}}} \right) - \frac{\nu_{22}}{\Sigma_{22}} \frac{\partial \sqrt{\Sigma_{33}}}{\partial u} &= 0 \end{aligned} \right\} \quad (I^*)$$

The relations (III) and (IV) which connect the coefficients of the two forms give the necessary and sufficient conditions which must be satisfied and we may say

Given the two differential quadratic forms

$$\begin{aligned} F &= \Sigma_{11} dt^2 + \Sigma_{22} du^2 + \Sigma_{33} dv^2 + 2\nu_{12} dt du + 2\nu_{13} dt dv + 2\nu_{23} du dv \\ \phi &= \nu_{11} dt^2 + \nu_{22} du^2 + \nu_{33} dv^2 + 2\nu_{12} dt du + 2\nu_{13} dt dv + 2\nu_{23} du dv \end{aligned}$$



of which the first is definite, then in order that there exist a hypersurface admitting these as first and second fundamental forms it is necessary and sufficient that relations (III) and (II) be satisfied. The corresponding hypersurface is determined & within translation\* in hyperspace.

Suppose the hypersurface whose existence and uniqueness we wish to demonstrate under the hypotheses above for referred to lines of curvature. Then (as will be shown later)  $E_{12} = E_{13} = E_{23} = D_{12} = D_{13} = D_{23} = 0$  and except in the case where  $E_{11} = 0$ , these lines of curvature are uniquely determined. Consider at every point of the hypersurface the characteristic tetrahedroid formed by the tangents to

\* Jordan - Essays sur la géométrie à  $n$  dimensions. Célébration de la 20<sup>ème</sup> année.



the positive directions of these lines  
and the normal to the hypersurface.  
Let  $(X_1, Y_1, Z_1, W_1)$ ,  $(X_2, Y_2, Z_2, W_2)$ ,  $(X_3, Y_3, Z_3, W_3)$ ,  
 $(X_4, Y_4, Z_4, W_4)$  be the respective direction co-  
sines of these lines

We then have

$$X_1 = \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial x}{\partial t}, \quad Y_1 = \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial y}{\partial t}, \quad Z_1 = \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial z}{\partial t}, \quad W_1 = \frac{1}{\sqrt{\epsilon_{11}}} \frac{\partial w}{\partial t}$$

$$X_2 = \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial x}{\partial t}, \quad \dots \dots \dots$$

$$X_3 = \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial x}{\partial t}, \quad \dots \dots \dots$$

$$X_4 = X \quad \dots \dots \dots$$

From formulae (I) and (II), page 19,  
substituting for the Christoffel symbols  
their values from page 15, we have

$$\frac{\partial X_1}{\partial t} = -\frac{X_2}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial u} - \frac{X_3}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial v} + \frac{\lambda_{11}}{\sqrt{\epsilon_{11}}} X_4$$

$$\frac{\partial X_1}{\partial u} = \frac{X_2}{\sqrt{\epsilon_{11}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial t}, \quad \frac{\partial X_1}{\partial v} = \frac{X_3}{\sqrt{\epsilon_{11}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial t}$$

$$\frac{\partial X_2}{\partial t} = \frac{X_1}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial u}, \quad \frac{\partial X_2}{\partial v} = \frac{X_3}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial u}$$

$$\frac{\partial X_2}{\partial u} = -\frac{X_1}{\sqrt{\epsilon_{11}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial t} - \frac{X_3}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial v} + \frac{\lambda_{22}}{\sqrt{\epsilon_{22}}} X_4$$



$$\frac{\partial X_3}{\partial t} = \frac{X_1}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial t}, \quad \frac{\partial X_3}{\partial u} = \frac{X_2}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial u}$$

$$\frac{\partial X_3}{\partial v} = -\frac{X_1}{\sqrt{\epsilon_{11}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial t} - \frac{X_2}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial u} + \frac{\partial \epsilon_{33}}{\partial \epsilon_{33}} X_4$$

$$\frac{\partial X_4}{\partial t} = -\frac{\partial_{11}}{\sqrt{\epsilon_{11}}} X_1, \quad \frac{\partial X_4}{\partial u} = -\frac{\partial_{22}}{\sqrt{\epsilon_{22}}} X_2, \quad \frac{\partial X_4}{\partial v} = -\frac{\partial_{33}}{\sqrt{\epsilon_{33}}} X_4$$

The unknown functions  $(X, X_1, X_2, X_3, X_4)$   
 $(Y, Y_1, Y_2, Y_3, Y_4), (Z, Z_1, Z_2, Z_3, Z_4), (W, W_1, W_2, W_3, W_4)$  then  
 satisfy the four simultaneous linear  
 homogeneous partial differential equations

$$\left. \begin{aligned} dQ_1 &= \left\{ -\frac{Q_2}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial u} - \frac{Q_3}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial v} + \frac{\partial_{11}}{\sqrt{\epsilon_{11}}} Q_1 \right\} dt + \frac{Q_2}{\sqrt{\epsilon_{11}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial t} du - \frac{Q_3}{\sqrt{\epsilon_{11}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial t} dv \\ dQ_2 &= \frac{Q_1}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial u} dt + \left\{ \frac{Q_1}{\sqrt{\epsilon_{11}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial t} - \frac{Q_3}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial v} + \frac{\partial_{22}}{\sqrt{\epsilon_{22}}} Q_2 \right\} du + \frac{Q_3}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial u} dv \\ dQ_3 &= \frac{Q_1}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial v} dt + \frac{Q_2}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{22}}}{\partial v} du + \left\{ \frac{Q_1}{\sqrt{\epsilon_{11}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial t} - \frac{Q_2}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{33}}}{\partial u} + \frac{\partial_{33}}{\sqrt{\epsilon_{33}}} Q_3 \right\} dv \\ dQ_4 &= -\frac{\partial_{11}}{\sqrt{\epsilon_{11}}} Q_1 dt - \frac{\partial_{22}}{\sqrt{\epsilon_{22}}} Q_2 du - \frac{\partial_{33}}{\sqrt{\epsilon_{33}}} Q_3 dv \end{aligned} \right\} \quad (4)$$

This is an infinitely integrable system, or consider the condition of integrability

$$\frac{\partial}{\partial u} \left( \frac{\partial Q_1}{\partial v} \right) = \frac{\partial}{\partial v} \left( \frac{\partial Q_1}{\partial u} \right)$$





Developing this it reduces to

$$\begin{aligned} & \left[ P_{11} \left( \frac{1}{\Sigma_{11}} \frac{\partial \Sigma_1}{\partial u} \right) + \frac{\partial}{\partial t} \left( \frac{1}{\Sigma_{11}} \frac{\partial \Sigma_{11}}{\partial t} \right) + \frac{1}{\Sigma_{33}} \frac{\partial \Sigma_{11}}{\partial r} \frac{\partial \Sigma_{11}}{\partial r} + \frac{\partial_{11} \partial_{11}}{\Sigma_{11} \Sigma_{11}} \right] \odot_1 \\ & + \left[ -P_{11} \left( \frac{1}{\Sigma_{33}} \frac{\partial \Sigma_{11}}{\partial r} \right) + \frac{1}{\Sigma_{11} \Sigma_{33}} \frac{\partial \Sigma_{11}}{\partial u} \frac{\partial \Sigma_{11}}{\partial r} \right] \odot_3 \\ & + \left[ P_{11} \left( \frac{\partial_{11}}{\Sigma_{11}} \right) - \frac{\partial_{11} \partial_{11}}{\Sigma_{11}} \frac{\partial \Sigma_{11}}{\partial u} \right] \odot_4 = 0 \end{aligned}$$

which is satisfied in virtue of equations (III\*\*) and (IV\*\*)

Since there exists an integral system and a single one, which for initial values of the variables

$$t = t_0, u = u_0, r = r_0$$

reduces to arbitrarily given initial values if  $(\odot_1, \odot_2, \odot_3, \odot_4), (\odot'_1, \odot'_2, \odot'_3, \odot'_4)$  are two integral systems then

$$\odot_1 \odot'_1 + \odot_2 \odot'_2 + \odot_3 \odot'_3 + \odot_4 \odot'_4 = \text{const.}$$

since the total differential of the first member is identically zero in virtue of equations (4).



Let  $(X, X_2, X_3, X_4), (Y, Y_2, Y_3, Y_4), (Z, Z_2, Z_3, Z_4), (W, W_2, W_3, W_4)$  be four integral systems  $\mathcal{O}(4)$  which for  $t = t_0, u = u_0, v = v_0$  reduce to the sixteenth coefficients of an orthogonal substitution. Then it follows from the observation above that for all values of the variables these sixteen quantities be the coefficients of an orthogonal substitution, and in particular

$$\bar{X}_i + \bar{Y}_i + \bar{Z}_i + \bar{W}_i = 1 \quad (i = 1, 2, 3, 4)$$

$$\bar{X}_i \bar{X}_j + \bar{Y}_i \bar{Y}_j + \bar{Z}_i \bar{Z}_j + \bar{W}_i \bar{W}_j = 0 \quad (i, j = 1, 2, 3, 4)$$

By (4) it is easily seen that the four expressions

$$\bar{E}_{11} \odot dt + \bar{E}_{21} \odot_2 du + \bar{E}_{31} \odot_3 dv, \quad \odot = X, Y, Z, W,$$

are exact differentials and writing

$$z = \int \bar{E}_{11} \odot_1 dt + \bar{E}_{21} \odot_2 du + \bar{E}_{31} \odot_3 dv$$

$$z = x, y, z, w, \quad \odot_i = X_i, Y_i, Z_i, W_i.$$

We have a hypersurface with the given fundamental forms.



Consider the system of equations (4).  
 This is identical with the system (34) (34')  
 (34') found by Professor Craig\*, where  
 $\alpha, \beta, \gamma$  are replaced by  $\mathcal{O}_1, \mathcal{O}_2, \mathcal{O}_3, \mathcal{O}_4$ , and  
 the  $\delta_{ij}$  by the coefficients in (4). Now  
 Professor Craig has shown that the  
 integration of system (34) can be re-  
 duced to the integration of a general-  
 ized\*\* Riccati equation and in the  
 same way the integration of system (4)  
 can be reduced to the integration of  
 the three simultaneous generalized  
 Riccati equations below by the  
 substitutions.

$$\mathcal{O}_1 = \frac{\gamma \lambda}{\lambda^2 + 1}, \quad \mathcal{O}_2 = \frac{\gamma u}{\lambda^2 + 1}, \quad \mathcal{O}_3 = \frac{\gamma v}{\lambda^2 + 1}, \quad \mathcal{O}_4 = \frac{\lambda^2 - 1}{\lambda^2 + 1}$$

$$\lambda^2 = \lambda^2 + u^2 + v^2$$

$$\frac{\partial \lambda}{\partial t} = - \left( \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial u} u - \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial v} v + \frac{\lambda^2 - 1}{2} \frac{\partial}{\partial \lambda} - \lambda^2 \frac{\partial}{\partial \lambda} \right)$$

$$\frac{\partial u}{\partial t} = \left( \frac{1}{\sqrt{\epsilon_{22}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial u} \lambda - \lambda u \frac{\partial}{\partial \lambda} \right)$$

$$\frac{\partial v}{\partial t} = \left( \frac{1}{\sqrt{\epsilon_{33}}} \frac{\partial \sqrt{\epsilon_{11}}}{\partial v} \lambda - \lambda v \frac{\partial}{\partial \lambda} \right)$$

\* American Journal of Math. Vol. XX No. 2, p. 45.

\*\* loc. cit., p. 41.



$$\left. \begin{aligned} \frac{\partial \lambda}{\partial u} &= \frac{1}{\sqrt{\varepsilon_{11}}} \frac{\partial \sqrt{\varepsilon_{11}}}{\partial t} u & -u \lambda \frac{\partial_{11}}{\sqrt{\varepsilon_{11}}} \\ \frac{\partial u}{\partial u} &= -\frac{1}{\sqrt{\varepsilon_{11}}} \frac{\partial \sqrt{\varepsilon_{11}}}{\partial t} \lambda - \frac{1}{\sqrt{\varepsilon_{33}}} \frac{\partial \sqrt{\varepsilon_{11}}}{\partial v} v + \frac{(K^2-1)\partial_{11}}{2\sqrt{\varepsilon_{12}}} & -u^2 \frac{\partial_{11}}{\sqrt{\varepsilon_{11}}} \\ \frac{\partial v}{\partial u} &= & \frac{1}{\sqrt{\varepsilon_{33}}} \frac{\partial \sqrt{\varepsilon_{11}}}{\partial v} u & -u \lambda \frac{\partial_{12}}{\sqrt{\varepsilon_{12}}} \end{aligned} \right\}$$

$$\left. \begin{aligned} \frac{\partial \lambda}{\partial v} &= \frac{1}{\sqrt{\varepsilon_{11}}} \frac{\partial \sqrt{\varepsilon_{33}}}{\partial t} v & -v \lambda \frac{\partial_{33}}{\sqrt{\varepsilon_{33}}} \\ \frac{\partial u}{\partial v} &= & -\frac{1}{\sqrt{\varepsilon_{11}}} \frac{\partial \sqrt{\varepsilon_{33}}}{\partial u} v & -u \frac{\partial_{33}}{\sqrt{\varepsilon_{33}}} \\ \frac{\partial v}{\partial v} &= -\frac{1}{\sqrt{\varepsilon_{11}}} \frac{\partial \sqrt{\varepsilon_{33}}}{\partial t} \lambda - \frac{1}{\sqrt{\varepsilon_{11}}} \frac{\partial \sqrt{\varepsilon_{33}}}{\partial u} u + \frac{(K^2-1)\partial_{33}}{2\sqrt{\varepsilon_{33}}} & -v^2 \frac{\partial_{33}}{\sqrt{\varepsilon_{33}}} \end{aligned} \right\}$$

Here we have chosen the special tetrahedroid formed by the tangents to the lines of curvature and the normal but any other tetrahedroidal tetrahedroid might have been chosen and we should have arrived at a set of equations similar to (4), illimitably integrable in virtue of relations (III) and (IV), and these would have led to the generalized Riccati equations of the form





$$\left. \begin{aligned} \frac{\partial \lambda}{\partial t} &= a_1 u + b_1 v + \left(\frac{\kappa^2 - 1}{2}\right) \lambda^2 + \lambda (c_1' \lambda + b_1' \mu + c_1' v) \\ \frac{\partial u}{\partial t} &= a_2 \lambda + b_2 v + \left(\frac{\kappa^2 - 1}{2}\right) \mu^2 + \mu (a_2' \lambda + b_2' \mu + c_2' v) \\ \frac{\partial v}{\partial t} &= a_3 \lambda + b_3 u + \left(\frac{\kappa^2 - 1}{2}\right) v^2 + v (c_3' \lambda + b_3' \mu + c_3' v) \end{aligned} \right\}$$

We may then say:

Given the two fundamental forms found, the first—being definite, and the coefficients—of which satisfy (III) and (IV), there exists a unique hypersurface admitting these as first and second fundamental forms, and I effectively obtain this hypersurface—it is necessary I integrate the simultaneous generalized Riccati equations.

The integration of a single generalized Riccati when solutular solutions are known has been studied by Mr. John Eiselein.\*

\* Note on the Integration of a Certain Class of Differential Equations—  
 American Journal of Math.—Vol. II No. 3.



### III

#### Lines on the Hypersurface

Extending to hypersurfaces the definition of lines of curvature there are lines along which the normals to the hypersurface form a developable surface. Let  $M$  be a point on the hypersurface and  $M_1$  the point where the normal touches the edge of regression of the developable.

Along a line of curvature  $x, y, z, w$  and  $X, Y, Z, W$  are functions of a single parameter the arc  $s$  say. Then  $x = X - rX, y = Y - rY, z = Z - rZ, w = W - rW$  where  $r$  is the algebraic value of  $MM_1$ .

Differentiating with respect to the arc.

$$\lambda X = \frac{dx}{ds} - r \frac{dX}{ds} - X \frac{dr}{ds}$$

$$\lambda Y = \frac{dy}{ds} - r \frac{dY}{ds} - Y \frac{dr}{ds}$$

$$\lambda Z = \frac{dz}{ds} - r \frac{dZ}{ds} - Z \frac{dr}{ds}$$

$$\lambda W = \frac{dw}{ds} - r \frac{dW}{ds} - W \frac{dr}{ds}$$



Multiplying by  $X, Y, Z, W$  in order and summing we have  $\lambda = -\frac{r}{S}$  and hence

$$\frac{dx}{S} = r \frac{\partial X}{\partial S}, \quad \frac{dy}{S} = r \frac{\partial Y}{\partial S}$$

$$\frac{dz}{S} = r \frac{\partial Z}{\partial S}, \quad \frac{dw}{S} = r \frac{\partial W}{\partial S}$$

Since along a line of curvature

$$\frac{dx}{\partial X} = \frac{dy}{\partial Y} = \frac{dz}{\partial Z} = \frac{dw}{\partial W} = r$$

Writing these four equations in curvilinear coordinates,

$$\frac{\partial x}{\partial t} dt + \frac{\partial x}{\partial u} du + \frac{\partial x}{\partial v} dv = r \left( \frac{\partial X}{\partial t} dt + \frac{\partial X}{\partial u} du + \frac{\partial X}{\partial v} dv \right)$$

$$\frac{\partial y}{\partial t} dt + \frac{\partial y}{\partial u} du + \frac{\partial y}{\partial v} dv = r \left( \frac{\partial Y}{\partial t} dt + \frac{\partial Y}{\partial u} du + \frac{\partial Y}{\partial v} dv \right)$$

$$\frac{\partial z}{\partial t} dt + \frac{\partial z}{\partial u} du + \frac{\partial z}{\partial v} dv = r \left( \frac{\partial Z}{\partial t} dt + \frac{\partial Z}{\partial u} du + \frac{\partial Z}{\partial v} dv \right)$$

$$\frac{\partial w}{\partial t} dt + \frac{\partial w}{\partial u} du + \frac{\partial w}{\partial v} dv = r \left( \frac{\partial W}{\partial t} dt + \frac{\partial W}{\partial u} du + \frac{\partial W}{\partial v} dv \right)$$

Multiplying by  $\frac{\partial t}{\partial x}, \frac{\partial t}{\partial y}, \frac{\partial t}{\partial z}, \frac{\partial t}{\partial w}, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial z}, \frac{\partial u}{\partial w}$  in order and summing each time we have the equivalent system



$$\left. \begin{aligned} E_{11} dt + E_{12} du + E_{13} dv &= -r (D_{11} dt + D_{12} du + D_{13} dv) \\ E_{21} dt + E_{22} du + E_{23} dv &= -r (D_{21} dt + D_{22} du + D_{23} dv) \\ E_{31} dt + E_{32} du + E_{33} dv &= -r (D_{31} dt + D_{32} du + D_{33} dv) \end{aligned} \right\} (1)$$

These are the equations of the lines of curvature on the hypersurface. Ricci\* has arrived at the same equations for the case of 3 variables by analogous reasoning and Foss\*\* has deduced similar ones, also for 4 variables, from considerations on the geodesic lines in a curved space. This last method shows at once that the lines of curvature are orthogonal and we may say.

Through every point on the hypersurface pass three lines of curvature forming a triply orthogonal system.

The three radii of curvature at the point are given by the roots of the cubic.  
\* Loc. Cit., p. 64.  
\*\* Loc. Cit., p. 40.





$$\begin{vmatrix} E_{11} + r D_{11} & E_{12} + r D_{12} & E_{13} + r D_{13} \\ E_{21} + r D_{21} & E_{22} + r D_{22} & E_{23} + r D_{23} \\ E_{31} + r D_{31} & E_{32} + r D_{32} & E_{33} + r D_{33} \end{vmatrix} = 0 \quad (2)$$

The product of the reciprocals of the three radii of curvature

$$\frac{1}{r_1 r_2 r_3} = - \frac{D}{4}$$

may be looked upon as the generalization of the total curvature of Gauss, and as will be shown in the section on spherical representations, the ratio of the element of hypersphere on the sphere to the element of hypersphere on the hypersurface is equal to the total curvature of the hypersurface.

Developing equation (2) above we have

$$\frac{1}{r^3} - \frac{H}{r^2} + \frac{L}{r} - K = 0$$

Palmer\* has given a rather compact form of the coefficients in this equation

\* *Sulle coordinate invariabili*  
*Memorie della R. Acc. dei Lincei Anno 1874*



Let  $T(y_1, y_2, y_3)$  be the equation of the hypersurface, and write

$$A_1 T = \frac{1}{2} \left( \frac{\partial T}{\partial y_1} \right)^2$$

$$A_{21} T = \frac{1}{2} \frac{\partial^2 T}{\partial y_1^2}$$

$$A_{22} T = \frac{1}{2} \left[ \frac{\partial^2 T}{\partial y_1^2} \frac{\partial^2 T}{\partial y_2^2} - \left( \frac{\partial^2 T}{\partial y_1 \partial y_2} \right)^2 \right]$$

From this deduce the following values for  $H$  and  $L$ .

$$L = \frac{1}{r_1} + \frac{1}{r_2} + \frac{1}{r_3} = \frac{A_{21} T}{\sqrt{A_1 T}} - \frac{A \sqrt{A_1 T}}{A T}$$

$$H = \frac{1}{r_1 r_2} + \frac{1}{r_1 r_3} + \frac{1}{r_2 r_3} = \frac{A_{22} T}{A_1 T} - \frac{A_{21} T}{\sqrt{A_1 T}} \cdot \frac{\sqrt{A_1 T}}{A T} + \frac{A_1 A_1 T}{4 (A_1 T)^2}$$

A particular case of this is of course, that of a surface immersed in three dimensional Euclidean space and Lamé\* has already given similar expressions for this. On a second \* because in les coordonnées curvilignes



hyper surface\* extends these base surface immersed in a hypersurface.

Writing  $T(x_1, x_2, x_3) = \text{const}$  as the equation of the surface immersed in the hypersurface whose linear element is given by

$$ds^2 = \sum_{i,j=1}^3 a_{ij} dx_i dx_j$$

he finds

$$\frac{1}{\rho_1} + \frac{1}{\rho_2} = \frac{4_1 U}{\sqrt{4_1 U}} - \frac{2\sqrt{4_1 U}}{4U}$$

$$\rho_1 \rho_2 = \frac{4_{12} U}{4U} - \frac{4_1 U}{\sqrt{4_1 U}} \frac{d\sqrt{4_1 U}}{dU} + \frac{d4_1 U}{4(4_1 U)^{3/2}}$$

where  $4_1 U$ ,  $4_{12} U$ ,  $4_{22} U$  are differential parameters constructed with respect to the linear element of the hypersurface and  $\rho_1, \rho_2$  satisfy

$$\begin{vmatrix} b_{11} - \rho \frac{4_{11}}{\sqrt{4_1 U}} & b_{12} - \rho \frac{4_{12}}{\sqrt{4_1 U}} \\ b_{12} - \rho \frac{4_{12}}{\sqrt{4_1 U}} & b_{22} - \rho \frac{4_{22}}{\sqrt{4_1 U}} \end{vmatrix} = 0$$



in which  $b_{rs} = \sum_{ij} a_{ij} \frac{\partial x_i}{\partial \xi_r} \frac{\partial x_j}{\partial \xi_s}$

$$\Omega_{rs} = \sum_{ik} \frac{\partial U}{\partial x_k} \frac{\partial x_k}{\partial \xi_r} \frac{\partial x_k}{\partial \xi_s} - \sum_{ik} \left[ \frac{\partial U}{\partial x_k} \right] \frac{\partial x_k}{\partial \xi_r} \frac{\partial x_k}{\partial \xi_s}$$

where  $\xi_1$  and  $\xi_2$  are curvilinear coordinates on  $T = \text{const.}$

We see that that the partial differential equations which define the minima and developable surfaces immersed in a hypersurface have the same form when expressed in the differential parameters as those for surfaces in Euclidean space of three dimensions.

If we suppose the hypersurface referred to a right orthogonal system  $x_1, x_2, x_3$ , then the condition that the parametric surface

$$x_1 = \text{const}$$

is a minimum surface is





$$\frac{\partial \log a_{12} a_{33}}{\partial x_1} = 0$$

the linear element of the hypersurface being

$$ds^2 = a_{11} dx_1^2 + a_{22} dx_2^2 + a_{33} dx_3^2$$

Let  $C$  be a curve on the hypersurface and consider  $t, u, v, x, y, z, w$  as functions of the arc. The direction cosines of the tangent  $OC$  at any point are given by

$$\cos \alpha = \frac{dx}{ds} = \frac{dx}{dt} \frac{dt}{ds} + \frac{dx}{du} \frac{du}{ds} + \frac{dx}{dv} \frac{dv}{ds}$$

$$\cos \beta = \frac{dy}{ds} = \frac{dy}{dt} \frac{dt}{ds} + \frac{dy}{du} \frac{du}{ds} + \frac{dy}{dv} \frac{dv}{ds}$$

$$\cos \gamma = \frac{dz}{ds} = \frac{dz}{dt} \frac{dt}{ds} + \frac{dz}{du} \frac{du}{ds} + \frac{dz}{dv} \frac{dv}{ds}$$

$$\cos \delta = \frac{dw}{ds} = \frac{dw}{dt} \frac{dt}{ds} + \frac{dw}{du} \frac{du}{ds} + \frac{dw}{dv} \frac{dv}{ds}$$

Let  $\sigma$  ( $0 < \sigma < \pi$ ) be the angle between the principal normal to  $C$  and the normal to the hypersurface. The direction cosines of the principal normal  $OC$  are\*

\* Piraudine Loc. cit. p. 22)



$$= \frac{d \cos \sigma}{ds} = \frac{d \cos \sigma}{ds} = \frac{d \cos \sigma}{ds} = \frac{d \cos \sigma}{ds}$$

where  $\rho$  is radius of first curvature of  $C$ .  
We then have

$$\cos \sigma = \rho \sum X \frac{d \cos \sigma}{ds}$$

Since we have at once by differentiating  
as indicated above

$$\begin{aligned} \frac{d \cos \sigma}{ds} &= \frac{d_1 dt' + d_2 du' + d_3 dv' + 2d_{12} dt du + 2d_{13} dt dv + 2d_{23} du dv}{\sqrt{E_1 dt'^2 + E_2 du'^2 + E_3 dv'^2 + 2E_{12} dt du + 2E_{13} dt dv + 2E_{23} du dv}} \\ &= \frac{d}{ds} \end{aligned}$$

Through the normal to the hyper-  
surface and the tangent to  $C$  pass  
a plane making a normal section.  
The first curvature of this section  
is given by the formula above,  
making  $\cos \sigma = \pm 1$ . Let  $R$  denote the  
radius of first curvature of this  
normal section, we then have

$$\frac{1}{R} = \pm \frac{\cos \sigma}{\rho}$$



Then  $\rho = \pm R \cos \theta$

Mensur's theorem then applies to the hypersurface, and therefore

The radius of first curvature of curve  $C$  traced on a hypersurface equals in every point  $M$  the radius of first curvature of the normal section made by the plane through the tangent to  $C$  at  $M$  multiplied by the cosine of the angle which this plane makes with the osculating plane of the curve  $C$ .

If  $R$  be taken positive when the section from the centre of curvature of the normal section toward the point  $M$  is along the positive normal to the hypersurface, then

$$\frac{1}{R} = -\frac{1}{\rho} \cos \theta$$



Refer the hypersurface dir-lines of curvature and denote the principal radii of curvature by  $r_1, r_2, r_3$ . Then along the  $t, u, v$  lines respectively

$$dx = r_1 dI, dy = r_1 dV, dz = r_1 dZ, dw = r_1 dW$$

$$dx = r_2 dI, \dots \dots \dots$$

$$dx = r_3 dI, \dots \dots \dots$$

or

$$\frac{\partial x}{\partial t} = r_1 \frac{\partial I}{\partial t}, \frac{\partial y}{\partial t} = r_1 \frac{\partial V}{\partial t}, \frac{\partial z}{\partial t} = r_1 \frac{\partial Z}{\partial t}, \frac{\partial w}{\partial t} = r_1 \frac{\partial W}{\partial t}$$

$$\frac{\partial x}{\partial u} = r_2 \frac{\partial I}{\partial u}, \dots \dots \dots$$

$$\frac{\partial x}{\partial v} = r_3 \frac{\partial I}{\partial v}, \dots \dots \dots$$

From these relations follow at once

$$E_{11} = -r_1 \rho_{11}, E_{22} = -r_2 \rho_{22}, E_{33} = -r_3 \rho_{33}$$

$$\rho_{12} = \rho_{13} = \rho_{23} = 0$$

hence

$$\frac{1}{R} = \frac{\frac{E_{11}}{r_1} dt^2 + \frac{E_{22}}{r_2} du^2 + \frac{E_{33}}{r_3} dv^2}{ds^2}$$

$$= \frac{E_{11}}{r_1} \left(\frac{dt}{ds}\right)^2 + \frac{E_{22}}{r_2} \left(\frac{du}{ds}\right)^2 + \frac{E_{33}}{r_3} \left(\frac{dv}{ds}\right)^2$$





Directed by  $\alpha, \beta, \gamma$  the direction angles  
of the tangent & this normal section  
referred to the lines of curvature,  

$$\frac{1}{R} = \frac{\cos^2 \alpha}{r_1} + \frac{\cos^2 \beta}{r_2} + \frac{\cos^2 \gamma}{r_3}$$

Since we know at once, that  $r_1, r_2, r_3$   
are the radii of first curvature of  
the normal sections tangent to  
the lines of curvature.

In the tangent hyperplane fix  
a system of rectangular axes  $\xi, \eta, \zeta$ ,  
coinciding with the positive directions  
of the lines of curvature  $\alpha, \beta, \gamma$  and  
consider the conicoid

$$\frac{\xi^2}{r_1} + \frac{\eta^2}{r_2} + \frac{\zeta^2}{r_3} = 1$$

The length of any semi-diameter  
whose direction angles are  $\alpha, \beta, \gamma$   
is given by  $\frac{1}{r} = \frac{\cos^2 \alpha}{r_1} + \frac{\cos^2 \beta}{r_2} + \frac{\cos^2 \gamma}{r_3}$



Since  $\sigma^2 = R$   
 and the square of every semidiameter  
 of the indicatrix equals the radius  
 of first curvature of the normal  
 section whose plane passes through  
 it

Define as conjugate directions  
 at a point on the hypersurface the  
 directions of three conjugate diam-  
 eters of the indicatrix at that  
 point. Let  $\alpha_1, \beta_1, \gamma_1; \alpha_2, \beta_2, \gamma_2; \alpha_3, \beta_3, \gamma_3$   
 denote the direction cosines of the  
 conjugate tangents. Then

$$\frac{\alpha_1 \alpha_2}{r_1} + \frac{\beta_1 \beta_2}{r_2} + \frac{\gamma_1 \gamma_2}{r_3} = 0$$

$$\frac{\alpha_1 \alpha_3}{r_1} + \frac{\beta_1 \beta_3}{r_2} + \frac{\gamma_1 \gamma_3}{r_3} = 0$$

$$\frac{\alpha_2 \alpha_3}{r_1} + \frac{\beta_2 \beta_3}{r_2} + \frac{\gamma_2 \gamma_3}{r_3} = 0$$

Still supposing the hypersurface re-  
 sulted did not have curvature, denote



displacements along the conjugate axes by  $u, v, w$ . Then

$$u = \sqrt{\epsilon_{11}} \frac{\delta t}{\delta s}, \quad v = \sqrt{\epsilon_{22}} \frac{\delta u}{\delta s}, \quad w = \sqrt{\epsilon_{33}} \frac{\delta v}{\delta s}$$

$$u_2 = \sqrt{\epsilon_{11}} \frac{\delta t}{\delta s}, \quad v_2 = \sqrt{\epsilon_{22}} \frac{\delta u}{\delta s}, \quad w_2 = \sqrt{\epsilon_{33}} \frac{\delta v}{\delta s}$$

$$u_3 = \sqrt{\epsilon_{11}} \frac{\delta t}{\delta s}, \quad v_3 = \sqrt{\epsilon_{22}} \frac{\delta u}{\delta s}, \quad w_3 = \sqrt{\epsilon_{33}} \frac{\delta v}{\delta s}$$

and the conjugate axes are given

$$u_1 \frac{\delta t}{\delta s} + \frac{\epsilon_{12}}{\epsilon_{11}} u_2 \frac{\delta u}{\delta s} + \frac{\epsilon_{13}}{\epsilon_{11}} u_3 \frac{\delta v}{\delta s} = 0$$

$$\frac{\epsilon_{11}}{\epsilon_{11}} \delta t \delta t + \frac{\epsilon_{12}}{\epsilon_{11}} \delta u \delta u + \frac{\epsilon_{13}}{\epsilon_{11}} \delta v \delta v = 0$$

$$\frac{\epsilon_{11}}{\epsilon_{11}} \delta t \delta t + \frac{\epsilon_{12}}{\epsilon_{11}} \delta u \delta u + \frac{\epsilon_{13}}{\epsilon_{11}} \delta v \delta v = 0$$

Now construct with respect to the second fundamental form, the equations analogous to the conditions of orthogonality for the lines constructed with respect to the first form. These three equations will give the conjugate lines in the hypersurface referred to any system



of coordinate lines. For we have

$$\begin{aligned} & \partial_1 dt \delta t + \partial_{12} du \delta u + \partial_{13} dv \delta v + \\ & \partial_{12} (\delta t \delta u + \delta u \delta t) + \partial_{13} (\delta t \delta v + \delta v \delta t) + \partial_{23} (\delta u \delta v + \delta v \delta u) = 0 \\ & \partial_{11} \delta t \delta t + \partial_{12} \delta u \delta u + \partial_{13} \delta v \delta v + \\ & \partial_{12} (\delta t \delta u + \delta u \delta t) + \partial_{13} (\delta t \delta v + \delta v \delta t) + \partial_{23} (\delta u \delta v + \delta v \delta u) = 0 \\ & \partial_{11} \delta t \delta t + \partial_{12} \delta u \delta u + \partial_{13} \delta v \delta v + \\ & \partial_{12} (\delta t \delta u + \delta u \delta t) + \partial_{13} (\delta t \delta v + \delta v \delta t) + \partial_{23} (\delta u \delta v + \delta v \delta u) = 0 \end{aligned}$$

and referring the hypersurface delta lines of curvature these reduce exactly to the three equations given above.

If the parametric lines form a conjugate system then these equations must be satisfied by

$$\begin{aligned} \delta t &= \delta t \quad \delta u = \delta v = 0 \\ \delta t &= 0, \quad \delta u = \delta u, \quad \delta v = 0 \\ \delta t &= 0, \quad \delta u = 0, \quad \delta v = \delta v \end{aligned}$$

hence  $\partial_{12} = \partial_{13} = \partial_{23} = 0$  and if these conditions are satisfied the equations above give the parametric lines.





Since the necessary and sufficient conditions that the coordinate lines form a conjugate system are

$$D_{11} = D_{22} = D_{33} = 0$$

Extending the definition of asymptotic lines on a surface a line is said to be asymptotic on a hypersurface if it coincides with its two conjugates. The asymptotic lines through a point then lie on the surface  $D_{11}dt^2 + D_{22}du^2 + D_{33}dv^2 + 2D_{12}dtdu + 2D_{13}dtdv + 2D_{23}du dv = 0$  and at the point this coincides with the asymptotic cone of the indicatrix. In order that the parametric lines lie on this surface we must have

$$D_{11} = D_{22} = D_{33} = 0$$

Since if the hypersurface is referred to any three orthogonal asymptotic lines  $D_{11} = D_{22} = D_{33} = 0$  and conversely,



If a plane moves so as to cut constantly tangent to a developable surface it generates a developable hypersurface. The developable to which the moving plane is constantly tangent is termed the surface of regression and the edge of regression of the developable surface is termed the edge of regression of the hypersurface. The rectilinear generatrices of the surface of regression are three lines of curvature on the hypersurface and  $K=0$ . Conversely if  $K=0$  we have a hypersurface with at least one line of curvature a right line. For determining lines of curvature we have

$$D_{11} D_{22} D_{33} = 0$$

and hence  $D_{11}, D_{22}, D_{33}$  must vanish.  
From (II) p. 19 we have



$$\frac{\partial X}{\partial t} = \frac{\partial Y}{\partial t} = \frac{\partial Z}{\partial t} = \frac{\partial W}{\partial t} = 0$$

from which it follows that  $X, Y, Z, W$ , are functions of  $u$  and  $v$  only. We also have

$$\frac{1}{E_{11}} \frac{\partial X}{\partial u} X + \frac{1}{E_{11}} \frac{\partial X}{\partial v} Y + \frac{1}{E_{11}} \frac{\partial X}{\partial u} Z + \frac{1}{E_{11}} \frac{\partial X}{\partial v} W = 0$$
$$\frac{1}{E_{11}} \frac{\partial X}{\partial u} \frac{1}{\partial u} + \dots = 0$$
$$\frac{1}{E_{11}} \frac{\partial X}{\partial v} \frac{\partial X}{\partial v} + \dots = 0$$

since the first expresses the orthogonality of the  $t$  line and the normal, while the second and third express the fact that  $d_{12} = d_{13} = 0$

The direction cosines

$$\frac{1}{E_{11}} \frac{\partial X}{\partial t}, \frac{1}{E_{11}} \frac{\partial Y}{\partial t}, \frac{1}{E_{11}} \frac{\partial Z}{\partial t}, \frac{1}{E_{11}} \frac{\partial W}{\partial t}$$

of the  $t$  line are thus functions of  $u$  and  $v$  only hence the  $t$  line is a right line and the hypersurface a line of zero curvature straight line.

The hypersurface of zero curvature is applicable to the hyperplane



For if  $D=0$  then the form  
 $\delta^2 = \epsilon_{11} du^2 + \epsilon_{22} dv^2 + \epsilon_{33} dw^2 + 2\epsilon_{12} du dv + 2\epsilon_{13} du dw + 2\epsilon_{23} dv dw$   
 is of class zero\*, and hence can be  
 deduced from a differential quad-  
 ratic form

$$\delta^2 = d\psi_1^2 + d\psi_2^2 + d\psi_3^2.$$

#### IV

On the Differential Equations satisfied  
by the Coordinates of a Point on the  
Hypersurface.

From equations (I), page 19 we  
 see that if  $\partial_1^2 = \partial_2^2 = \partial_3^2 = 0$  the coor-  
 dinates of any point on the hypersur-  
 face satisfy the simultaneous partial  
 differential equations of the form

$$\left. \begin{aligned} \frac{\partial^2 \theta}{\partial x_1 \partial x_2} &= a_1 \frac{\partial \theta}{\partial x_1} + b_1 \frac{\partial \theta}{\partial x_2} + c_1 \frac{\partial \theta}{\partial x_3} \\ \frac{\partial^2 \theta}{\partial x_1 \partial x_3} &= a_2 \frac{\partial \theta}{\partial x_1} + b_2 \frac{\partial \theta}{\partial x_2} + c_2 \frac{\partial \theta}{\partial x_3} \\ \frac{\partial^2 \theta}{\partial x_2 \partial x_3} &= a_3 \frac{\partial \theta}{\partial x_1} + b_3 \frac{\partial \theta}{\partial x_2} + c_3 \frac{\partial \theta}{\partial x_3} \end{aligned} \right\}$$

+

\* see - loc. cit., p. 56.





Conversely, if the four simultaneous equations above admit four common solutions  $q_1, q_2, q_3, q_4$ , linearly independent, then the formulae

$q_1 = q_1(x_1, x_2, x_3), q_2 = q_2(x_1, x_2, x_3), q_3 = q_3(x_1, x_2, x_3), q_4 = q_4(x_1, x_2, x_3)$  define a hypersurface referred to conjugate lines.

For

$$\left| \begin{array}{cccc} \frac{\partial q_1}{\partial x_1 \partial x_2} & \frac{\partial q_1}{\partial x_1 \partial x_3} & \frac{\partial q_2}{\partial x_2 \partial x_3} & \frac{\partial q_4}{\partial x_1 \partial x_3} \\ \frac{\partial q_1}{\partial x_2} & - & - & - \\ \frac{\partial q_2}{\partial x_2} & - & - & - \\ \frac{\partial q_3}{\partial x_2} & - & - & - \\ \frac{\partial q_4}{\partial x_3} & - & - & - \end{array} \right| \begin{array}{l} (i, j = 1, 2, 3) \\ = 0 \\ (i \neq j) \end{array}$$

and hence  $\Delta_1 = \Delta_2 = \Delta_3 = 0$

If in addition  $q_1, q_2, q_3, q_4$  which  $\Delta = q_1^2 + q_2^2 + q_3^2 + q_4^2$  are common solutions then

$$\frac{\partial^2 \Delta}{\partial x_1 \partial x_2} \frac{\{1, 1\}}{\partial x_1} - \frac{\{1, 2\}}{\partial x_1} \frac{\partial q_1}{\partial x_2} - \frac{\{1, 3\}}{\partial x_1} \frac{\partial q_2}{\partial x_3} = 2\Delta_{11} = 0$$

Hence  $\Delta_1 = \Delta_2 = \Delta_3 = 0$  and the formulae above define a hypersurface referred



& its line of curvature.

Again if  $\rho_{11} = \rho_{22} = \rho_{33} = 0$  the coordinates satisfy the simultaneous equations of the type

$$\left. \begin{aligned} \frac{\partial^2 \phi}{\partial x_1^2} &= a_1 \frac{\partial \phi}{\partial x_1} + b_1 \frac{\partial \phi}{\partial x_2} + c_1 \frac{\partial \phi}{\partial x_3} \\ \frac{\partial^2 \phi}{\partial x_2^2} &= a_2 \frac{\partial \phi}{\partial x_1} + b_2 \frac{\partial \phi}{\partial x_2} + c_2 \frac{\partial \phi}{\partial x_3} \\ \frac{\partial^2 \phi}{\partial x_3^2} &= a_3 \frac{\partial \phi}{\partial x_1} + b_3 \frac{\partial \phi}{\partial x_2} + c_3 \frac{\partial \phi}{\partial x_3} \end{aligned} \right\} \quad (B)$$

and conversely if these three equations admit four simultaneous solutions  $\phi_1, \phi_2, \phi_3, \phi_4$  linearly independent then the formulae

$$\phi_1 = \phi(x_1, x_2), \quad \phi_2 = \phi(x_1, x_3), \quad \phi_3 = \phi(x_2, x_3), \quad \phi_4 = \phi(x_1, x_2, x_3)$$

define a hypersurface referred to three asymptotic lines.

For

$$\begin{vmatrix} \frac{\partial^2 \phi_1}{\partial x_1^2} & \frac{\partial^2 \phi_1}{\partial x_2^2} & \frac{\partial^2 \phi_1}{\partial x_3^2} & \frac{\partial^2 \phi_1}{\partial x_1^2} \\ \frac{\partial^2 \phi_2}{\partial x_1^2} & \frac{\partial^2 \phi_2}{\partial x_2^2} & \frac{\partial^2 \phi_2}{\partial x_3^2} & \frac{\partial^2 \phi_2}{\partial x_1^2} \\ \frac{\partial^2 \phi_3}{\partial x_1^2} & \frac{\partial^2 \phi_3}{\partial x_2^2} & \frac{\partial^2 \phi_3}{\partial x_3^2} & \frac{\partial^2 \phi_3}{\partial x_1^2} \\ \frac{\partial^2 \phi_4}{\partial x_1^2} & \frac{\partial^2 \phi_4}{\partial x_2^2} & \frac{\partial^2 \phi_4}{\partial x_3^2} & \frac{\partial^2 \phi_4}{\partial x_1^2} \end{vmatrix} = 0 \quad (C)$$



and hence

$$\mathcal{D}_{11} = \mathcal{D}_{22} = \mathcal{D}_{33} = 0$$

A projective transformation preserves the conjugate and asymptotic lines for a projective transformation is defined by

$$\eta_1 = \frac{\alpha_1}{\epsilon}, \eta_2 = \frac{\alpha_2}{\epsilon}, \eta_3 = \frac{\alpha_3}{\epsilon}, \eta_4 = \frac{\alpha_4}{\epsilon}$$

where  $\alpha_1, \alpha_2, \alpha_3, \alpha_4$  are linear functions

of  $\eta_1, \eta_2, \eta_3, \eta_4$  and  $\epsilon$  is a constant.

The coordinate lines form a conjugate system if  $\alpha_1, \alpha_2, \alpha_3, \alpha_4$  are solutions of (A) and if an asymptotic system, (B).

Substituting

$$\mathcal{D}' = \frac{\alpha}{\epsilon}$$

(A) is transformed into a system of the same kind in the first case, and the same is true of (B) in the second case. The theorem is then proved.

Inversion preserves the lines of curvature. This transformation is given by



$$y_1 = \frac{\kappa^2 y_1'}{\rho}, y_2 = \frac{\kappa^2 y_2'}{\rho}, y_3 = \frac{\kappa^2 y_3'}{\rho}, y_4' = \frac{\kappa^2 y_4'}{\rho}$$

where  $\rho = y_1'^2 + y_2'^2 + y_3'^2 + y_4'^2$

Refer the hypersurface  $S$  to the curvatures  $\rho, \sigma, \tau$ . Then  $y_1, y_2, y_3, y_4, \rho$  satisfy the three simultaneous equations

$$\frac{\partial^2 \rho}{\partial y_i \partial y_j} = a_{ij} \frac{\partial \rho}{\partial y_1} + b_{ij} \frac{\partial \rho}{\partial y_2} + c_{ij} \frac{\partial \rho}{\partial y_3} \quad (i, j = 1, 2, 3) \quad (C)$$

Now  $y_1, y_2, y_3, y_4$  can be expressed in terms of  $y_1, y_2, y_3, y_4$  as follows

$$y_1 = \frac{\kappa^2 y_1'}{\rho}, y_2 = \frac{\kappa^2 y_2'}{\rho}, y_3 = \frac{\kappa^2 y_3'}{\rho}, y_4 = \frac{\kappa^2 y_4'}{\rho}$$

where  $\sigma = y_1'^2 + y_2'^2 + y_3'^2 + y_4'^2$  and  $\tau = \frac{\kappa^4}{\rho^2}$

Effect an equations (C), the substitution

$$\rho = \frac{\sigma}{\tau}$$

The equations in  $\tau$  admit as particular solutions  $y_1, y_2, y_3, y_4$  and 1, and are therefore of the form

$$\frac{\partial^2 \tau}{\partial y_i \partial y_j} = a_{ij} \frac{\partial \tau}{\partial y_1} + b_{ij} \frac{\partial \tau}{\partial y_2} + c_{ij} \frac{\partial \tau}{\partial y_3} \quad (i, j = 1, 2, 3) \quad (C')$$

Now (C) admits the solution  $\sigma = 1$





and hence (B) admits the solution  $\infty$ .  
 The hypersurface  $\infty$  at the point  
 $M(\frac{1}{4}, \frac{1}{4}, \frac{1}{4}, \frac{1}{4})$  has  $\kappa_1, \kappa_2, \kappa_3$  as its  
 principal curvatures and these are therefore preserved.

Consider now the three simultaneous  
 equations

$$\frac{\partial^2 \theta}{\partial x_1 \partial x_2} = 0, \quad \frac{\partial^2 \theta}{\partial x_1 \partial x_3} = 0, \quad \frac{\partial^2 \theta}{\partial x_2 \partial x_3} = 0$$

The general integral is of the form

$$\theta = f(x_1) + \phi(x_2) + \psi(x_3)$$

where  $f, \phi, \psi$  are arbitrary functions.

Writing

$$\theta_i = f_i(x_1) + \phi_i(x_2) + \psi_i(x_3) \quad i=1,2,3,4$$

we have a hypersurface of translation referred to a conjugate system.  
 This may be generated by moving any one of the parametric  
 surfaces so as to have its points  
 describe curves contained in the  
 intersection of the other two.



The differential equation of the surface on which the asymptotic lines through a point lie when the hypersurface is referred to a conjugate system, becomes

$$\partial_{11} dx_1^2 + \partial_{22} dx_2^2 + \partial_{33} dx_3^2 = 0$$

and for a hypersurface of translation this takes the form,

$$\begin{vmatrix} f & f' & f'' & f''' \\ f' & f'' & f''' & f^{(4)} \\ f'' & f''' & f^{(4)} & f^{(5)} \\ f''' & f^{(4)} & f^{(5)} & f^{(6)} \end{vmatrix} dx_1^2 + \begin{vmatrix} \phi & \phi' & \phi'' & \phi''' \\ \phi' & \phi'' & \phi''' & \phi^{(4)} \\ \phi'' & \phi''' & \phi^{(4)} & \phi^{(5)} \\ \phi''' & \phi^{(4)} & \phi^{(5)} & \phi^{(6)} \end{vmatrix} dx_2^2 + \begin{vmatrix} \psi & \psi' & \psi'' & \psi''' \\ \psi' & \psi'' & \psi''' & \psi^{(4)} \\ \psi'' & \psi''' & \psi^{(4)} & \psi^{(5)} \\ \psi''' & \psi^{(4)} & \psi^{(5)} & \psi^{(6)} \end{vmatrix} dx_3^2 = 0$$

If now  $f = f' = f'' = f''' = \phi' = \phi'' = 0$

this reduces at once to

$$\begin{vmatrix} \psi & \psi' \\ \psi' & \psi'' \end{vmatrix} dx_1^2 + \begin{vmatrix} \phi & \phi' \\ \phi' & \phi'' \end{vmatrix} dx_2^2 + \begin{vmatrix} \psi & \psi' \\ \psi' & \psi'' \end{vmatrix} dx_3^2 = 0$$

in which the variables can obviously be separated. Hence

The asymptotic lines are contained by quadratics in all hypersurfaces of



translation whose generating lines lie in three mutually orthogonal hyperplanes.

The equations of the normal at any point on the hypersurface are

$$\sum (y-y') \frac{\partial y}{\partial x_1} = 0 \quad \sum (y-y') \frac{\partial y}{\partial x_2} = 0 \quad \sum (y-y') \frac{\partial y}{\partial x_3} = 0$$

or

$$\left. \begin{aligned} \sum y' \frac{\partial y}{\partial x_1} - \frac{\partial x}{\partial x_1} &= 0 \\ \sum y' \frac{\partial y}{\partial x_2} - \frac{\partial x}{\partial x_2} &= 0 \\ \sum y' \frac{\partial y}{\partial x_3} - \frac{\partial x}{\partial x_3} &= 0 \end{aligned} \right\} \quad (1)$$

where  $r = \sum \frac{x^2}{2}$

By imposing the condition that the normal have unit a developable surface we have for a point  $y'$ , suitably chosen on the normal

$$\sum \frac{\partial y}{\partial x_1} dy' = 0 \quad \sum \frac{\partial y}{\partial x_2} dy' = 0 \quad \sum \frac{\partial y}{\partial x_3} dy' = 0$$

Differentiating (1) under this hypothesis

$$\left. \begin{aligned} \sum y' d \frac{\partial y}{\partial x_1} - d \frac{\partial x}{\partial x_1} &= 0 \\ \sum y' d \frac{\partial y}{\partial x_2} - d \frac{\partial x}{\partial x_2} &= 0 \\ \sum y' d \frac{\partial y}{\partial x_3} - d \frac{\partial x}{\partial x_3} &= 0 \end{aligned} \right\} \quad (2)$$



Eliminating  $y_1, y_2, y_3, y_4$  between (1) and (2) we have the three equations

$$\begin{vmatrix} \lambda \frac{\partial^2 \phi}{\partial x_1^2} & \lambda \frac{\partial^2 \phi}{\partial x_1 \partial x_2} & \frac{\partial^2 \phi}{\partial x_1 \partial x_3} & \frac{\partial^2 \phi}{\partial x_1 \partial x_4} & \frac{\partial^2 \phi}{\partial x_1 \partial x_5} \\ \lambda \frac{\partial^2 \phi}{\partial x_2^2} & - & - & - & - \\ \lambda \frac{\partial^2 \phi}{\partial x_3^2} & - & - & - & - \\ \lambda \frac{\partial^2 \phi}{\partial x_4^2} & - & - & - & - \\ \lambda \frac{\partial^2 \phi}{\partial x_5^2} & - & - & - & - \end{vmatrix} = 0 \quad \begin{matrix} i, j = 1, 2, 3 \\ i \neq j \end{matrix}$$

which give the lines of curvature.

Consider now the three simultaneous equations

$$A_{ij} \frac{\partial^2 \phi}{\partial x_i^2} + B_{ij}' \frac{\partial^2 \phi}{\partial x_i \partial x_j} + C_{ij} \frac{\partial^2 \phi}{\partial x_j^2} + A_{ji} \frac{\partial^2 \phi}{\partial x_j^2} + B_{ji}' \frac{\partial^2 \phi}{\partial x_j \partial x_i} + C_{ji} \frac{\partial^2 \phi}{\partial x_i^2} = 0 \quad (4)$$

$i \neq j, \quad i, j = 1, 2, 3,$

and write the conditions that these admit the given particular solutions  $y_1, y_2, y_3, y_4, \dots$ . Equations (4) can then be written in the determinant form

$$\begin{vmatrix} \frac{\partial^2 \phi}{\partial x_1^2} & \frac{\partial^2 \phi}{\partial x_1 \partial x_2} & \frac{\partial^2 \phi}{\partial x_1 \partial x_3} & \frac{\partial^2 \phi}{\partial x_1 \partial x_4} & \frac{\partial^2 \phi}{\partial x_1 \partial x_5} \\ \frac{\partial^2 \phi}{\partial x_2^2} & - & - & - & - \\ \frac{\partial^2 \phi}{\partial x_3^2} & - & - & - & - \\ \frac{\partial^2 \phi}{\partial x_4^2} & - & - & - & - \\ \frac{\partial^2 \phi}{\partial x_5^2} & - & - & - & - \end{vmatrix} = 0$$





are using equations (3) with (3a) in that case can be written

$$C_{ij} \dot{x}_i^2 - C_{ij} \dot{x}_i \dot{x}_j - A_{ij} \dot{x}_i^2 + L_{ij} \dot{x}_i \dot{x}_j + M_{ij} \dot{x}_i \dot{x}_j + N_{ij} \dot{x}_i^2 = 0 \quad (6)$$

where  $K = 1, 2, 3, i \neq j \neq K$

where

$$L_{ij} = \left| \frac{\partial u}{\partial x_i^2} \frac{\partial u}{\partial x_j \partial x_K} \right| + \left| \frac{\partial u}{\partial x_i \partial x_K} \frac{\partial u}{\partial x_j \partial x_i} \right|$$

$$M_{ij} = \left| \frac{\partial u}{\partial x_i \partial x_j} \frac{\partial u}{\partial x_i \partial x_K} \right| + \left| \frac{\partial u}{\partial x_i \partial x_K} \frac{\partial u}{\partial x_j^2} \right|$$

$$N_{ij} = \left| \frac{\partial u}{\partial x_i \partial x_K} \frac{\partial u}{\partial x_j \partial x_K} \right|$$

$$\left| \frac{\partial u}{\partial x_i \partial x_j} \frac{\partial u}{\partial x_i \partial x_K} \right| = \begin{vmatrix} \frac{\partial u}{\partial x_i \partial x_j} & \frac{\partial u}{\partial x_i \partial x_K} & \frac{\partial u}{\partial x_j} & \frac{\partial u}{\partial x_K} & \frac{\partial u}{\partial x_i} \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ \frac{\partial u}{\partial x_i \partial x_j} & \frac{\partial u}{\partial x_i \partial x_K} & \frac{\partial u}{\partial x_j} & \frac{\partial u}{\partial x_K} & \frac{\partial u}{\partial x_i} \end{vmatrix}$$

Now  $\left| \frac{\partial u}{\partial x_i \partial x_j} \frac{\partial u}{\partial x_i \partial x_K} \right|$  can at once be written as  $\left| \frac{\partial u}{\partial x_i^2} \frac{\partial u}{\partial x_j^2} \right|, \left| \frac{\partial u}{\partial x_i^2} \frac{\partial u}{\partial x_j^2} \right|, \left| \frac{\partial u}{\partial x_i^2} \frac{\partial u}{\partial x_j^2} \right|$ ,  $\frac{\partial u}{\partial x_i^2} \frac{\partial u}{\partial x_j^2}$ ,  $\frac{\partial u}{\partial x_i^2} \frac{\partial u}{\partial x_j^2}$  by means of equation (3).



since  $L_{ij}$ ,  $M_{ij}$ ,  $N_{ij}$  are linearly ex-  
pressible in terms of  $\gamma_{ij}$ ,  $\beta_{ij}$ ,  $\alpha_{ij}$ .  
We see then that by any means we  
can obtain three partial differ-  
ential equations of the form (4) simul-  
taneously satisfied by  $\gamma$ ,  $\beta$ ,  $\alpha$ .  
Then the differential equations of the  
lines of curvature are given by (6).

In particular  $\gamma_{ij} = \beta_{ij} = \alpha_{ij} = 0$   
and obviously  $L_{ij} = M_{ij} = N_{ij} = 0$   
and (6) reduces to

$$\gamma_1 dx_1 = 0 \quad \beta_1 dx_1 = 0 \quad \alpha_1 dx_1 = 0$$

the hypersurface being referred to the  
lines of curvature as we have already  
seen. The transformation variables  
defined by (6) must then reduce  
 $\gamma$  and  $\beta$  simultaneously to zero and  
(6) plays the same role with regard  
to (4) that the equation of character-  
istics does with regard to the



single equation

$$+ \frac{12x}{x^2} + 6 \frac{y^2}{y^2} + 6 \frac{z^2}{z^2} + 4 \frac{12x}{x^2} + 15 \frac{y^2}{y^2} = 0$$

It may then say:

Moreover we have obtained the partial differential equation of the form (4) which are simultaneously satisfied by  $x, y, z$ . The differential equations of the conjugate lines are given by (6) and in addition (4) clearly  $x, y, z$  is a particular solution. It defines the lines of curvature.

We can now extend Dupin's definition conjugate lines of the hypersurface.

Let the tangent <sup>hyper</sup> plane at  $(x_1, y_1, z_1)$

$$x_1 - x_2 = p_1(x_1 - x_2) + p_2(y_1 - y_2) + p_3(z_1 - z_2) \quad (1)$$

where  $p_i = \frac{\partial \phi}{\partial x_i}$

suppose this tangent hyper plane displaced along the three coordinate lines  $x_1, x_2, x_3$ .



The characteristic in each case is a plane and these three planes are given by (V) and each of

$$\sum \frac{p}{x_1} (y-y_1) = 0 \quad \sum \frac{p}{x_2} (y-y_2) = 0 \quad \sum \frac{p}{x_3} (y-y_3) = 0 \quad (8)$$

Express now the conditions that the lines formed by taking these three planes in pairs are tangent to the coordinate lines. For this it is necessary and sufficient write that the ~~three~~ <sup>three</sup> functions formed by (7) and each of (8) are satisfied when  $(y_1 - y_1), (y_2 - y_2), (y_3 - y_3), (y_4 - y_4)$  are replaced by  $\frac{\partial y_1}{\partial x_1}, \frac{\partial y_2}{\partial x_2}, \frac{\partial y_3}{\partial x_3}, \frac{\partial y_4}{\partial x_4}$  properly substituted.

This leads at once to

$$\frac{\partial y_4}{\partial x_1 \partial x_2} - \rho_1 \frac{\partial y_4}{\partial x_1 \partial x_3} - \rho_2 \frac{\partial y_4}{\partial x_2 \partial x_3} - \rho_3 \frac{\partial y_4}{\partial x_2 \partial x_4} = 0 \quad (9)$$

(1, 2) = (2, 1) (1, 3) = (3, 1) (1, 4) = (4, 1) (2, 4) = (4, 2)

We have also

$$\frac{\partial y_4}{\partial x_1} = \rho_1 \frac{\partial y_4}{\partial x_2} + \rho_2 \frac{\partial y_4}{\partial x_3} + \rho_3 \frac{\partial y_4}{\partial x_4}$$





Eliminating  $\rho_1, \rho_2, \rho_3$  between (8) and (9) taken in pairs we get as the required conditions

$$\begin{vmatrix} \frac{\partial^2 \varphi_1}{\partial x_i \partial x_j} & \frac{\partial \varphi_1}{\partial x_i} & \frac{\partial \varphi_1}{\partial x_j} & \frac{\partial \varphi_1}{\partial x_3} \\ \frac{\partial^2 \varphi_2}{\partial x_i \partial x_j} & - & - & - \\ \frac{\partial^2 \varphi_3}{\partial x_i \partial x_j} & - & - & - \\ \frac{\partial^2 \varphi_4}{\partial x_i \partial x_j} & - & - & - \end{vmatrix} \quad \begin{matrix} i, j = 1, 2, 3 \\ i \neq j \end{matrix} = 0$$

These are the conditions that the point structural coordinates of the hypersurface satisfy three simultaneous partial differential equations of the type

$$\frac{\partial^2 \varphi}{\partial x_i \partial x_j} = a_{ij} \frac{\partial^2 \varphi}{\partial x_1^2} + a_{ij} \frac{\partial^2 \varphi}{\partial x_2^2} + c_{ij} \frac{\partial^2 \varphi}{\partial x_3^2} \quad (12)$$

(i, j = 1, 2, 3, i \neq j)

Thus we may define as conjugate lines on the hypersurface lines such that if the tangent



hyperplane is displaced along any line to the intersection of the characteristics is tangent to the line.

The equations satisfied by the tangential coordinates can also be easily derived by following the reasoning of Darboux\* in the case of the two dimensional.

Writing  $y_1, y_2, y_3, y_4, y_5$  the homogeneous coordinates of a point, the tangent hyperplane is

$$a_1 y_1' + a_2 y_2' + a_3 y_3' + a_4 y_4' + a_5 y_5' = 0 \quad (13)$$

where  $a_i$  is a function of  $y_1, y_2, y_3$  the independent parameters. The hypersurface is obtained by eliminating  $y_1, y_2, y_3$  between (13) and its three partial derivatives

$$\frac{\partial}{\partial y_1'} \frac{\partial}{\partial y_1'} = 0 \quad \frac{\partial}{\partial y_2'} \frac{\partial}{\partial y_2'} = 0 \quad \frac{\partial}{\partial y_3'} \frac{\partial}{\partial y_3'} = 0 \quad (14)$$

\* These results are found in



Since  $M(x, y, z) = 1$  the point  $S$  lies on the hyper-tangent plane, then  
 $\sum a_i = 0, \sum y_i \frac{\partial a_i}{\partial x_1} = 0, \sum y_i \frac{\partial a_i}{\partial x_2} = 0, \sum y_i \frac{\partial a_i}{\partial x_3} = 0$  (15)  
 which lead at once to

$$\sum a_i \frac{\partial^2 y_i}{\partial x_j^2} = 0, \sum a_i y_i \frac{\partial^2 y_i}{\partial x_j^2} = 0, \sum a_i y_i \frac{\partial^2 y_i}{\partial x_j^2} = 0 \quad (16)$$

Putting now the conditions that  $y_1, y_2, y_3$  form a conjugate system we have at once

$$\sum a_i y_i \frac{\partial^2 y_i}{\partial x_j^2} = 0, \sum y_i \frac{\partial^2 a_i}{\partial x_j^2} = 0 \quad (i, j = 1, 2, 3) \quad (17)$$

Eliminating  $y_i$  between (15) and the first three of (17) we get the equations satisfied by the tangential coordinates

$$\begin{vmatrix} \frac{\partial^2 a_1}{\partial x_1^2} & \frac{\partial a_1}{\partial x_1} & \frac{\partial a_1}{\partial x_2} & \frac{\partial a_1}{\partial x_3} & a_1 \\ \vdots & \vdots & \vdots & \vdots & \vdots \\ \frac{\partial^2 a_j}{\partial x_i^2} & \frac{\partial a_j}{\partial x_i} & \frac{\partial a_j}{\partial x_2} & \frac{\partial a_j}{\partial x_3} & a_j \end{vmatrix} = 0 \quad (18)$$

with the similar ones for the point coordinates. Hence



Given the three partial differential equations

$$C_1 \frac{\partial^2 z}{\partial x^2} + C_2 \frac{\partial^2 z}{\partial y^2} + C_3 \frac{\partial^2 z}{\partial x \partial y} - C_4 \frac{\partial z}{\partial x} - C_5 \frac{\partial z}{\partial y} + C_6 z = 0 \quad (17)$$

which are satisfied by the homogeneous means spherical or tangential coordinates of a hyperboloid the partial equations

$$C_1 \frac{\partial^2 z}{\partial x^2} + C_2 \frac{\partial^2 z}{\partial y^2} + C_3 \frac{\partial^2 z}{\partial x \partial y} - C_4 \frac{\partial z}{\partial x} - C_5 \frac{\partial z}{\partial y} + C_6 z = 0 \quad (20)$$

gives the conjugate lines, and it is easily seen that if (19) admits  $\frac{\partial^2 z}{\partial x^2} + \frac{\partial^2 z}{\partial y^2} + \frac{\partial^2 z}{\partial x \partial y}$  is a particular solution in the first case and  $\frac{\partial^2 z}{\partial x^2} + \frac{\partial^2 z}{\partial y^2} + \frac{\partial^2 z}{\partial x \partial y}$  in the second then so gives the linear generators.

# V

On the Spherical Representation of the general representation of Gauss, we suppose given a curve from the center of the hyperboloid of two sheets





about the origin as centre parallel to the normal to the hypersurface along any line there will meet the hypersphere the spherical representative of the line on the hypersurface. Write the corresponding linear element on the hypersphere

$$ds^2 = \xi^2 dt^2 + \xi_2^2 du^2 + \xi_3^2 dv^2 + 2\xi_2 \xi_3 dt du + 2\xi_3 \xi_1 dt dv + 2\xi_1 \xi_2 du dv \quad (1)$$

We have then at once from (II) page 19

$$\begin{aligned} \xi_1 &= \sqrt{\frac{X}{Y}} = \frac{1}{2} (\partial_1 |\partial_1'| + \partial_{21} |\partial_1''| + \partial_{31} |\partial_1'''|) \\ \xi_2 &= \sqrt{\frac{X}{Y}} = \frac{1}{2} (\partial_2 |\partial_1'| + \partial_{22} |\partial_1''| + \partial_{32} |\partial_1'''|) \\ \xi_{23} &= \sqrt{\frac{X}{Y}} = \frac{1}{2} (\partial_3 |\partial_1'| + \partial_{23} |\partial_1''| + \partial_{33} |\partial_1'''|) \\ \xi_{12} &= \sqrt{\frac{X}{Y}} \sqrt{\frac{X}{Y}} = \frac{1}{2} (\partial_2 |\partial_1'| + \partial_{22} |\partial_1''| + \partial_{32} |\partial_1'''|) \\ &= \frac{1}{2} (\partial_1 |\partial_1'| + \partial_{21} |\partial_1''| + \partial_{31} |\partial_1'''|) \\ \xi_3 &= \sqrt{\frac{X}{Y}} \sqrt{\frac{X}{Y}} = \frac{1}{2} (\partial_3 |\partial_1'| + \partial_{23} |\partial_1''| + \partial_{33} |\partial_1'''|) \\ &= \frac{1}{2} (\partial_1 |\partial_1'| + \partial_{21} |\partial_1''| + \partial_{31} |\partial_1'''|) \\ \xi_{23} &= \sqrt{\frac{X}{Y}} \sqrt{\frac{X}{Y}} = \frac{1}{2} (\partial_3 |\partial_1'| + \partial_{23} |\partial_1''| + \partial_{33} |\partial_1'''|) \\ &= \frac{1}{2} (\partial_1 |\partial_1'| + \partial_{21} |\partial_1''| + \partial_{31} |\partial_1'''|) \end{aligned} \quad (2)$$



From these it is easily seen that the truly orthogonal system formed by the lines of curvature on the hypersurface gives a truly orthogonal system on the hypersphere, not vice versa.

$$\varepsilon_{12} = \varepsilon_{13} = \varepsilon_{23} = \alpha_{12} = \alpha_{13} = \alpha_{23} = 0$$

we have at once

$$e_{12} = e_{13} = e_{23} = 0$$

Let us now see what relations exist between the angles made by a set of conjugate lines on the hypersurface and those made by their spherical representations.

In this case

$$\alpha_{12} = \alpha_{13} = \alpha_{23} = 0$$

$$e_{11} = \frac{\partial_{11}^2}{4} \left| \frac{\varepsilon_{11} \varepsilon_{33}}{\varepsilon_{32} \varepsilon_{33}} \right| \equiv \frac{\partial_{11}^2}{4} \varepsilon'_{11} \text{ say.}$$

$$e_{22} = \frac{\partial_{22}^2}{4} \varepsilon'_{22}, \quad e_{33} = \frac{\partial_{33}^2}{4} \varepsilon'_{33}$$

$$e_{12} = \frac{\partial_{12} \partial_{21}}{4} \varepsilon'_{12}, \quad e_{13} = \frac{\partial_{11} \partial_{33}}{4} \varepsilon'_{13}, \quad e_{23} = \frac{\partial_{22} \partial_{33}}{4} \varepsilon'_{23}$$



Let  $\Omega_{ij}$  are the same significance in the hyperplane as  $\omega_{ij}$  in the hyper surface. Then

$$\cos \Omega_{12} = \frac{e_{12}}{E_{11} E_{22}} = \frac{E'_{12}}{\sqrt{E_{11} E_{22}}}$$

$$\text{From } E_{12} = \sqrt{E_{11} E_{22}} \cos \omega_{12}$$

$$E_{13} = \sqrt{E_{11} E_{33}} \cos \omega_{13}$$

$$E_{23} = \sqrt{E_{22} E_{33}} \cos \omega_{23}$$

we have at once

$$E'_{12} = \left| \begin{array}{cc} \sqrt{E_{11} E_{22}} \cos \omega_{12} & \sqrt{E_{11} E_{33}} \cos \omega_{13} \\ \sqrt{E_{22} E_{33}} \cos \omega_{23} & E_{33} \end{array} \right| = E_{33} \sqrt{E_{11} E_{22}} \left| \begin{array}{cc} \cos \omega_{12} \cos \omega_{23} \\ \cos \omega_{13} & 1 \end{array} \right|$$

$$E'_{11} = \left| \begin{array}{cc} 1 & \cos \omega_{13} \\ \cos \omega_{13} & 1 \end{array} \right| E_{11} E_{33} = E_{11} E_{33} \sin^2 \omega_{13}$$

which give

$$\cos \Omega_{12} = \frac{\cos \omega_{12} - \cos \omega_{13} \cos \omega_{23}}{\sin \omega_{13} \sin \omega_{23}}$$

$$\cos \Omega_{13} = \frac{\cos \omega_{13} - \cos \omega_{12} \cos \omega_{23}}{\sin \omega_{12} \sin \omega_{23}}$$

$$\cos \Omega_{23} = \frac{\cos \omega_{23} - \cos \omega_{12} \cos \omega_{13}}{\sin \omega_{12} \sin \omega_{13}}$$



If two of the conjugate lines is taken in a surface of curvature the third is orthogonal to them, hence

The angle between two conjugate lines lying in a surface of curvature or a hypersurface is preserved or changed into its supplement by spherical representation.

Let  $\omega$  be the solid angle formed by any three formative surfaces at a point  $M$  on the hypersurface. The element of hyperarea at this point

$$dS = \sin \omega \, ds_1 \, ds_2 \, ds_3$$

and the corresponding element of hyperarea on the hypersphere is

$$dS' = \sin \Omega \, d\theta_1 \, d\theta_2 \, d\theta_3$$

It is easily seen that

$$\frac{dS'}{dS} = K$$

the total curvature of the hypersurface.





So here

$$\sin \tilde{\omega} = \begin{vmatrix} 1 & \cos \omega_{1,2} & \cos \omega_{1,3} \\ \cos \omega_{1,2} & 1 & \cos \omega_{2,3} \\ \cos \omega_{1,3} & \cos \omega_{2,3} & 1 \end{vmatrix} = \mathcal{E}_{11} \mathcal{E}_{33} \mathcal{I}$$

$$\sin \tilde{\Omega} = \begin{vmatrix} 1 & \cos \Omega_{1,2} & \cos \Omega_{1,3} \\ \cos \Omega_{1,2} & 1 & \cos \Omega_{2,3} \\ \cos \Omega_{1,3} & \cos \Omega_{2,3} & 1 \end{vmatrix} = \mathcal{E}_{11} \mathcal{E}_{22} \mathcal{E}_{33} \mathcal{I}$$

here  $\mathcal{I} = \begin{vmatrix} \mathcal{E}_{11} & \mathcal{E}_{12} & \mathcal{E}_{13} \\ \mathcal{E}_{12} & \mathcal{E}_{22} & \mathcal{E}_{23} \\ \mathcal{E}_{13} & \mathcal{E}_{32} & \mathcal{E}_{33} \end{vmatrix}$

Again  $\mathcal{E}_{11} \mathcal{E}_{22} \mathcal{E}_{33} = \frac{1}{\sqrt{\mathcal{E}_{11} \mathcal{E}_{22} \mathcal{E}_{33}}}$  it is so

$\mathcal{E}_{11} \mathcal{E}_{22} \mathcal{E}_{33}' = \frac{1}{\sqrt{\mathcal{E}_{11} \mathcal{E}_{22} \mathcal{E}_{33}}}$  it is so

hence  $\frac{\mathcal{E}_{33}'}{\mathcal{E}_{33}} = \frac{\mathcal{I}'}{\mathcal{I}}$

now it is easily seen, that

$$\sqrt{\mathcal{I}} = \begin{vmatrix} \frac{\partial \mathcal{X}}{\partial t} & \frac{\partial \mathcal{Y}}{\partial t} & \frac{\partial \mathcal{Z}}{\partial t} & \frac{\partial \mathcal{W}}{\partial t} \\ \frac{\partial \mathcal{X}}{\partial u} & - & - & - \\ \frac{\partial \mathcal{X}}{\partial v} & - & - & - \\ \frac{\partial \mathcal{X}}{\partial \lambda} & - & - & - \end{vmatrix} \quad \sqrt{\mathcal{I}'} = \begin{vmatrix} \frac{\partial \mathcal{X}}{\partial t} & \frac{\partial \mathcal{Y}}{\partial t} & \frac{\partial \mathcal{Z}}{\partial t} & \frac{\partial \mathcal{W}}{\partial t} \\ \frac{\partial \mathcal{X}}{\partial u} & - & - & - \\ \frac{\partial \mathcal{X}}{\partial v} & - & - & - \\ \frac{\partial \mathcal{X}}{\partial \lambda} & - & - & - \end{vmatrix}$$



From which it follows that we have

$$D = -\sqrt{S_1}$$

$$\text{or } K = -\frac{D}{4} = \sqrt{\frac{S_1}{4}}$$

$$\text{Hence } \frac{dS'}{dS} = K$$

We have in fact:

The limit of the ratio of the element of hypersphere to the hypersphere of the corresponding element of hypersphere on the hypersurface is equal to the total curvature of the hypersurface.

Consider now the following notation. Given the third differential quadratic form

$$ds'^2 = \sum_{i,j} c_{ij} dx_i dx_j \quad (3)$$

Let the  $c_{ij}$  be assigned functions of  $x_1, x_2, x_3$  required to correspond to the hypersurface.

Applying equations (I) and (II)



page 19. The hyperplanes we have

$$\frac{\partial \mathcal{Q}}{\partial x_i \partial x_j} = \sum_k \left\{ \begin{smallmatrix} i, j \\ k \end{smallmatrix} \right\} \frac{\partial \mathcal{Q}}{\partial x_k} - \ell_{ij} \mathcal{Q} \quad (i, j = 1, 2, 3, 4) \quad (4)$$

(i ≠ j)

where  $\left\{ \begin{smallmatrix} i, j \\ k \end{smallmatrix} \right\}$  is constructed with respect to (3).

Now excluding the hypersurface of zero curvature,  $\alpha_1, \alpha_2, \alpha_3, \alpha_4$  can be so determined that the four equations

$$\alpha_i = \alpha_1 \frac{\partial \mathcal{Q}_i}{\partial x_1} + \alpha_2 \frac{\partial \mathcal{Q}_i}{\partial x_2} + \alpha_3 \frac{\partial \mathcal{Q}_i}{\partial x_3} + \alpha_4 \mathcal{Q}_i \quad (i = 1, 2, 3, 4) \quad (5)$$

shall be satisfied, where  $\alpha_1, \alpha_2, \alpha_3, \alpha_4$  are any four functions of  $x_1, x_2, x_3$ , since the determinant of the system does not vanish, being equal to the square root of the discriminant of (3). Putting  $\alpha_i = \frac{\partial \mathcal{Q}_i}{\partial x_k}$  and multiplying the four equations (5) in order by  $\frac{\partial \mathcal{Q}_1}{\partial x_k}, \frac{\partial \mathcal{Q}_2}{\partial x_k}, \frac{\partial \mathcal{Q}_3}{\partial x_k}, \mathcal{Q}_k$  and summing in each case, we have

$$-\mathcal{D}_{K1} = \sum_j \alpha_j \ell_{Kj} \quad 0 = \alpha_4 \quad (K = 1, 2, 3)$$



Following these

$$\sigma_k = -\frac{1}{f} |D_k^*|' \quad (k = \dots, 3)$$

where in  $|D_k^*|'$   $E_{im}$  is replaced by  $E_{km}$

We have that it once the equations

$$\frac{\partial \Phi}{\partial x_i} = -\frac{1}{f} (|D_i^*|' \frac{\partial \Phi}{\partial x_1} + |D_i^*|' \frac{\partial \Phi}{\partial x_2} + |D_i^*|' \frac{\partial \Phi}{\partial x_3}) \quad (6)$$

( $i = 1, 2, 3$ )

These three equations are of exactly the same form as (II) page 7 with  $\Phi$  written for  $\phi$  and  $E_{ij}$  for  $E_{ij}$  and therefore the conditions of integrability for (6) are constructed with respect to (3) in exactly the same way that the conditions of integrability for (II) are constructed with respect to the same

$$ds^2 = \sum_{i,j} E_{ij} dx_i dx_j$$

Referring then to equations (II) page 7 we have it once so the conditions of integrability of (6),





$$\frac{\partial D_s}{\partial x_t} - \frac{\partial D_t}{\partial x_s} + \frac{3}{2} \left( D_s \{x_s'\} - D_t \{x_t'\} \right) \quad (2)$$

$(x_s, t = 1, 2, 3; \quad t \neq s)$

hence we may say:

Given the two differential quad-  
rics forms

$$Q = \sum_{s=1}^3 D_s dx_s$$
$$Q' = \sum_{s=1}^3 C_s dx_s$$

where the first is definite and  
signature +1 (belongs to the hyperbolic).  
Then in order that there exist a  
hypersurface admitting these as  
second and third fundamental  
forms, it is necessary and suf-  
ficient that equations (1) be sat-  
isfied and the corresponding  
expression is obtained by the  
integration of (6).







15

Vita

George Oscar Lumer was born  
near Hagerstown, Virginia August 1 1883.  
He received his elementary education  
at Churchland Academy near Hager-  
stown and entered the Johns Hopkins  
University in October 1892 as can-  
didate for the degree of Bachelor  
of Arts in the mathematical-  
physical group. In June 1895  
he was graduated and appointed  
University Scholar in Mathematics,  
selecting Physics and Astronomy  
as first and second subordi-  
nates. In October 1896 he  
resigned the appointment of Hon-  
orary Hopkins Scholar becoming  
Instructor in Physics in the  
University of Illinois and began



entered the Johns Hopkins  
University in October of 1897. In  
June of 1898 he was made  
Fellow in Mathematics.















